

轴向基础窄带随机激励悬臂梁的稳定性

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摘要 基于 Kane 方程, 建立了含耦合三次几何及惯性非线性项轴向基础激励悬臂梁动力学方程. 采用多尺度法, 对所得方程进行一次近似展开, 着重计算了窄带随机主参激共振时系统平凡响应的最大 Lyapunov 指数及随机稳定性问题, 并通过直接数值积分验证了所得稳定区域的有效性.

关键词 悬臂梁, 主参激共振, 窄带随机激励, 最大 Lyapunov 指数, 随机稳定性

引言

受基础激励的弹性结构在工程中随处可见. 几十年来, 许多学者在该领域, 特别是确定性基础激励悬臂梁的非线性动力学的研究, 取得了一系列的成果, 但在随机激励下悬臂梁的非线性动力学的研究方面鲜有报道.

Nayfeh 与 Pai 对轴向非伸长细长梁非线性参激响应问题进行了详尽的研究, 其结果显示, 梁的第一阶模态的主参激共振幅频响应曲线呈现硬特性, 而二阶及以上模态却呈软特性^[1]. Anderson, Nayfeh 与 Balachandran 对轴向基础简谐激励悬臂梁的动力响应进行了实验测定, 研究结果表明: 实验结果与理论预测间具有较好的一致性^[2].

在另一研究领域, Kane, Ryan 与 Banerjee^[3]采用 Kane 方程并结合假设模态, 系统地建立了非惯性场中梁的动力学方程. Yoo, Ryan 与 Scott^[4]对文献[3]的建模方法中个别变量进行了更规范化的选取, 使得建模结果更趋简单. Hyun 和 Yoo^[5]在文献[4]的基础上详细研究了受轴向基础简谐激励悬臂梁平凡响应稳定性问题等. 由于文献[3~5]在建模过程中对广义惯性力及广义作用力进行了线性化处理, 从而使对梁更复杂的非线性动力行为的研究无法进行. 最近, 冯志华、胡海岩采用文献[3~5]的建模方法, 在保留广义惯性力及广义作用力非线性项的基础上, 得出了大范围直线运动梁的动力学运动方程, 以悬臂梁为分析对象进行非线性动力学研究,

并同相关研究成果进行了比较, 从一个侧面验证了建模的正确性^[6], 同时, 涉及了主参激共振与内共振联合作用下梁的非线性动力行为研究^[7~8]、窄带随机参激激励梁的稳定性与随机分叉问题^[9]等.

本文基于文[6]的建模理论, 利用多尺度法对方程进行了一次近似展开, 着重研究了在窄带随机主参激共振时系统平凡响应的随机稳定性问题, 并进行了直接数值积分验证, 取得了较好的一致性.

1 运动方程的建立

图 1 为受轴向基础 A 激励悬臂梁模型. 根据文献[6]的建模方法, 最终梁的纵横振动方程为

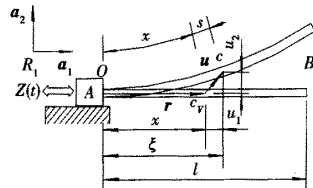


图 1 受轴向基础参激激励悬臂梁

Fig. 1 Configuration of a cantilever beam undergoing the axial excitation of its base

$$M_{1k} \ddot{q}_k + K_{1k} q_k - \sum_{i=1}^{n_2} \sum_{j=1}^{n_2} (2D_{ij}^k + E_{ij}^k) Q_i Q_j - \sum_{i=1}^{n_2} \sum_{j=1}^{n_2} C_{ij}^k (\dot{Q}_i \dot{Q}_j + Q_i \ddot{Q}_j) = -a_{1k} \ddot{z} \quad (1)$$

$$k = 1, 2, \dots, n_1$$

$$M_{2k} \ddot{Q}_k + K_{2k} Q_k - \sum_{i=1}^{n_2} a_{2i}^k \ddot{z}_i Q_i - \sum_{i=1}^{n_1} \sum_{j=1}^{n_2} C_{2ij}^k \ddot{q}_i Q_j -$$

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$$\begin{aligned} & \sum_{i=1}^{n_1} \sum_{j=1}^{n_2} (4E_{2ij}^{k1} + E_{2ij}^{k2} + E_{2ij}^{k3}) q_i \dot{Q}_j + \\ & \sum_{h=1}^{n_2} \sum_{i=1}^{n_2} \sum_{j=1}^{n_2} D_{2hij}^k (Q_h Q_i \ddot{Q}_j + Q_h \dot{Q}_i \dot{Q}_j) + \\ & \sum_{h=1}^{n_2} \sum_{i=1}^{n_2} \sum_{j=1}^{n_2} (F_{ijh}^k + G_{2ijh}^k) Q_h Q_i \dot{Q}_j = 0 \end{aligned} \quad (2)$$

$k = 1, 2, \dots, n$

式中相关参数见文[6].

2 一阶近似展开

设 $\ddot{z} = f \cos \Omega t + g \sin \Omega t$, 其中 Ω 为中心频率, f, g 分别是独立的慢时变高斯随机过程, 可由下式确定

$$\begin{cases} \dot{f} = (\gamma/2)^{1/2} W_c - (\gamma/2) f \\ \dot{g} = (\gamma/2)^{1/2} W_s - (\gamma/2) g \end{cases} \quad (3)$$

式中 γ 为窄带随机噪声带宽, W_c, W_s 分别为独立的高斯白噪声且有相同谱密度常数 S_0 . 假设 $z(t)$ 为小幅随机激励, 并设 Ω 远小于梁的第一阶纵向振动固有频率, 假定梁为细长梁, 参照文献[1, 2, 6~9]的处理方式, 在下述分析中认为梁是轴向非伸缩梁, 即不计及梁的纵向振动. 为此, 引入无量纲变量

$$\eta = \frac{x}{l}, \tau = 1.8715^2 \frac{t}{T}, y_k = \frac{Q_k}{l}, T = \sqrt{\frac{\rho l^4}{EI}} \quad (4)$$

将式(3)和式(4)与上述假设一并代入式(2), 并引入线性阻尼, 最终有

$$\begin{aligned} & \ddot{y}_k + 2\epsilon^2 \zeta_k \omega_k \dot{y}_k + \omega_k^2 y_k - \\ & \epsilon^2 \sum_{i=1}^{n_2} \tilde{z} \alpha_{2i}^k y_i + \sum_{i=1}^{n_2} \sum_{j=1}^{n_2} \sum_{h=1}^{n_2} \alpha_{ijh}^k y_i y_j y_h + \\ & \sum_{i=1}^{n_2} \sum_{j=1}^{n_2} \sum_{h=1}^{n_2} \beta_{ijh}^k y_i (\dot{y}_j \dot{y}_h + y_j \ddot{y}_h) = 0 \end{aligned} \quad (5)$$

$k = 1, 2, \dots, n_2$

式中 $\omega_k^2 = \frac{\bar{K}_{2k}}{\bar{M}_{2k}}, \bar{M}_{2k} = \int_0^1 \Phi_{2k}^2 d\eta, \omega = \frac{\Omega T}{1.8715^2},$

$\bar{K}_{2k} = \int_0^1 \Phi_{2k}^2 \epsilon^2 l \bar{M}_{2k} d\eta, \tilde{z} = \bar{f} \cos \omega \tau + \bar{g} \sin \omega \tau,$

$\bar{f} = \frac{f T^2}{1.8715^4 \epsilon^2 l \bar{M}_{2k}}, \bar{g} = \frac{g T^2}{1.8715^4 \epsilon^2 l \bar{M}_{2k}},$

$\tilde{\alpha}_{2i}^k = \int_0^1 \phi_{ik} d\eta, \alpha_{ijh}^k = \frac{\tilde{\alpha}_{ijh}^k}{\bar{M}_{2k}}, \beta_{ijh}^k = \frac{\tilde{\beta}_{ijh}^k}{\bar{M}_{2k}},$

$\tilde{\alpha}_{ijh}^k = \frac{1}{1.8715^4} \int_0^1 \Phi_{2i, \eta} \Phi_{2j, \eta} (\Phi_{2h, \eta} \Phi_{2k, \eta} + \Phi_{2h, \eta} \Phi_{2k, \eta}) d\eta,$

$\tilde{\beta}_{ijh}^k = \int_0^1 \phi_{ij} \phi_{hk} d\eta, \phi_{ij} = \int_0^{\eta} \Phi_{2i, \beta} \Phi_{2j, \beta} d\beta, \epsilon$ 为正的小参

数. 其中 $\tilde{\alpha}_{2i}^k, \alpha_{ijh}^k, \beta_{ijh}^k$ 及 \bar{f}, \bar{g} 为参激激励系数、非线性几何系数、非线性惯性系数及无量纲随机参激激励幅值.

将悬臂梁振型函数

$$\begin{aligned} \Phi_{2k}(\eta) &= \frac{1}{2} [(\cosh \lambda_k \eta - \cos \lambda_k \eta) - \\ & \frac{\cos \lambda_k + \cosh \lambda_k}{\sin \lambda_k + \sinh \lambda_k} (\sinh \lambda_k \eta - \sin \lambda_k \eta)] \end{aligned} \quad (6)$$

代入上述各系数式中即可求得相应的系数值, 而 $\bar{M}_{2k} = 0.25$.

采用多尺度法, 将

$$y_k(\tau, \epsilon) = \epsilon y_{k0}(T_0, T_2) + \epsilon^3 y_{k1}(T_0, T_2)$$

代入式(5), 对其进行一次近似展开并令 ϵ^1 及 ϵ^3 前系数为零, 相应地求得

$$y_{k0} = A_k(T_0) \exp(i\omega_k T_0) + cc, \quad k = 1, 2, \dots, n_2 \quad (7)$$

式中 cc 表示前项共轭(下同). 而 y_{k1} 则可由下列方程组决定

$$\begin{aligned} D_0^2 y_{k1} + \omega_k^2 y_{k1} &= -2i\omega_k (\zeta_k \omega_k A_k + A_k') \exp(i\omega_k T_0) + \\ & \frac{1}{2} \sum_{m=1}^{n_2} (\bar{f} - i\bar{g}) \tilde{\alpha}_{2m}^k \{A_m \exp[i(\omega + \omega_m) T_0] + \\ & \bar{A}_m \exp[i(\omega - \omega_m) T_0]\} + \\ & \sum_{m=1}^{n_2} \sum_{j=1}^{n_2} \sum_{h=1}^{n_2} \{(-\alpha_{mjh}^k + \omega_j \omega_h \beta_{mjh}^k + \omega_h^2 \beta_{mjh}^k) \times \\ & A_m A_j A_h \exp[i(\omega_m + \omega_j + \omega_h) T_0] + \\ & (-\alpha_{mjh}^k - \omega_j \omega_h \beta_{mjh}^k + \omega_h^2 \beta_{mjh}^k) \times \\ & A_m \bar{A}_j A_h \exp[i(\omega_m - \omega_j + \omega_h) T_0] + \\ & (-\alpha_{mjh}^k + \omega_j \omega_h \beta_{mjh}^k + \omega_h^2 \beta_{mjh}^k) \times \\ & \bar{A}_m A_j A_h \exp[i(-\omega_m + \omega_j + \omega_h) T_0] + \\ & (-\alpha_{mjh}^k - \omega_j \omega_h \beta_{mjh}^k + \omega_h^2 \beta_{mjh}^k) \times \\ & \bar{A}_m \bar{A}_j A_h \exp[i(-\omega_m - \omega_j + \omega_h) T_0]\} + cc \end{aligned} \quad (8)$$

$k = 1, 2, \dots, n_2$

当研究第 p 阶模态主参激共振情况时, 引入频率调谐因子 σ , 使

$$\omega = 2\omega_p + \epsilon^2 \sigma \quad p = 1, 2, \dots, n_2 \quad (9)$$

将式(9)代入式(8), 消除永年项. 由于阻尼的存在, 最终只有梁的第 p 阶模态对系统的长期动力行为起作用, 故有

$$\begin{aligned} 2i\omega_p (\zeta_p \omega_p A_p + A_p') - \frac{1}{2} (\bar{f} - i\bar{g}) \times \\ \tilde{\alpha}_{2p}^p \bar{A}_p \exp(i\sigma T_2) + c_p A_p^2 \bar{A}_p = 0 \end{aligned} \quad (10)$$

式中 $c_p = 3\alpha_{ppp}^p - 2\omega_p^2 \beta_{ppp}^p$ 为非线性系数.

3 主参激共振时系统平凡响应的稳定性

将式

$$A_p(T_2) = \frac{1}{2}[U_p(T_2) - iV_p(T_2)]\exp\left(\frac{i\sigma T_2}{2}\right) \quad (11)$$

代入式(10),最终有

$$\begin{cases} U_p = -\left(\zeta_p\omega_p + \frac{\bar{g}\bar{a}_{2p}^1}{4\omega_p}\right)U_p - \frac{1}{2}\left(\sigma - \frac{1}{2\omega_p}\mathcal{F}\bar{a}_{2p}^1\right)V_p + \frac{1}{8\omega_p}c_pV_p(U_p^2 + V_p^2) \\ V_p = \frac{1}{2}\left(\sigma + \frac{1}{2\omega_p}\mathcal{F}\bar{a}_{2p}^1\right)U_p + \left(-\zeta_p\omega_p + \frac{1}{4\omega_p}\bar{g}\bar{a}_{2p}^1\right)V_p - \frac{1}{8\omega_p}c_pU_p(U_p^2 + V_p^2) \end{cases} \quad (12)$$

不失一般性,现研究系统在第一阶模态主参激共振时平凡响应在概率一意义下的稳定性问题. 根据 Oseledec 乘法遍历性定理,上述系统平凡响应在概率一意义下的渐近稳定性的充分必要条件是系统线性化方程的最大 Lyapunov 指数小于零. 而线性化方程的 Lyapunov 指数定义如下

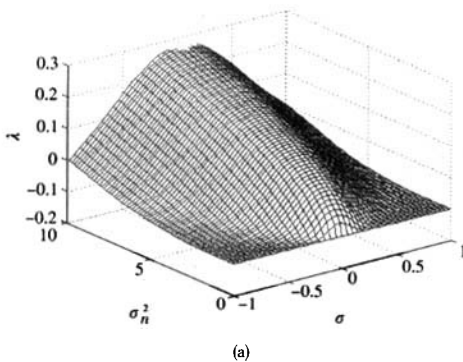
$$\lambda = \lim_{t \rightarrow \infty} \frac{1}{t} \ln \|Z(t; z_0)\| \quad (13)$$

式中 $Z(t; z_0)$ 为方程的解向量, z_0 为初始状态向量.

方程(12)在 $(U_1, V_1) = (0, 0)$ 邻域的线性化展开及相应的 Lyapunov 指数表达式分别为

$$\begin{cases} U_1 = -\left(\zeta_1\omega_1 + \frac{\bar{g}\bar{a}_{21}^1}{4\omega_1}\right)U_1 - \frac{1}{2}\left(\sigma - \frac{1}{2\omega_1}\mathcal{F}\bar{a}_{21}^1\right)V_1 \\ V_1 = \frac{1}{2}\left(\sigma + \frac{1}{2\omega_1}\mathcal{F}\bar{a}_{21}^1\right)U_1 + \left(-\zeta_1\omega_1 + \frac{1}{4\omega_1}\bar{g}\bar{a}_{21}^1\right)V_1 \end{cases} \quad (14)$$

及



$$\lambda = \lim_{T_2 \rightarrow \infty} \frac{1}{T_2} \ln \sqrt{(U_1(T_2, U_{10}, V_{10}))^2 + (V_1(T_2, U_{10}, V_{10}))^2} \quad (15)$$

根据文[9,10],在数值计算中,随机信号 $\tilde{z}(t)$ 采用下列形式

$$\tilde{z}(t) = \sqrt{\frac{2\sigma_n^2}{N}} \sum_{k=1}^N \cos(\omega_k t + \varphi_k) \quad (16)$$

式中 ω_k, φ_k 分别为 $(\omega - \gamma/2, \omega + \gamma/2)$ 及 $(0, 2\pi)$ 内均布随机变量序列, γ 为带宽. 根据相关文献可知: 当 $N \rightarrow \infty$ 时, $E[\tilde{z}(t)] = \sigma_n^2$ 且 $\tilde{z}(t)$ 趋于 Gauss 随机过程. 在实际计算时,结合文献[9,10],取 $N = 500$,有

$$\begin{cases} \mathcal{F}(T_2) = \sqrt{\frac{2\sigma_n^2}{N}} \sum_{k=1}^N \cos\left(\frac{\gamma}{\epsilon^2}\Omega_k T_2 + \varphi_k\right) \\ \bar{g}(T_2) = -\sqrt{\frac{2\sigma_n^2}{N}} \sum_{k=1}^N \sin\left(\frac{\gamma}{\epsilon^2}\Omega_k T_2 + \varphi_k\right) \end{cases} \quad (17)$$

式中 Ω_k 为 $(-0.5, 0.5)$ 内均布随机变量序列. 经计算, $\omega_1 = 1.0000, \bar{a}_{21}^1 = 0.3927$.

图 2 为 $\gamma=0.01, \zeta_1=0.1$ 时系统平凡响应的最大 Lyapunov 指数及相应的等高线图 ($\lambda > 0$ 区域为系统平凡响应的几乎确定不稳定区域). 图 3 为 $\gamma=0.01, \zeta_1=0.2$ 时系统平凡响应的最大 Lyapunov 指数及相应的等高线图,从图中可见,阻尼比增大,不稳定区域明显变小. 为了检验结果的正确与否,现直接对式(14)进行数值积分计算,计算中选择: $(\sigma_n^2, \sigma) = (2, 0), (y_1(0), \dot{y}_1(0)) = (0, 1)$, 根据所给条件, $(\sigma_n^2, \sigma) = (2, 0)$ 处在图 2(b)的不稳定区域,其结果见图 4(a), $(\sigma_n^2, \sigma) = (2, 0)$ 处在图 3(b)的稳定区域,

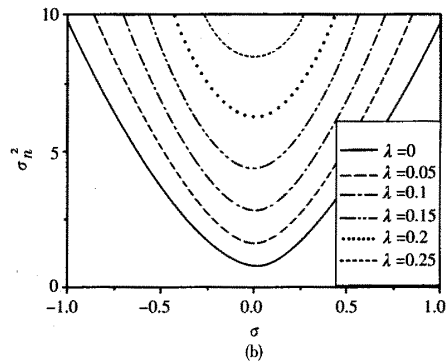


图 2 最大 Lyapunov 指数及等高线图 ($\gamma=0.01, \zeta_1=0.1$)

Fig. 2 The largest Lyapunov exponents and the corresponding isohypses

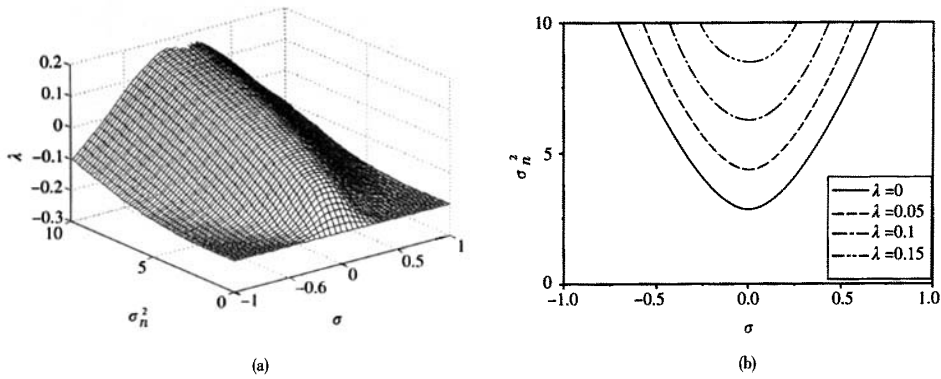


图 3 最大 Lyapunov 指数及等高线图($\gamma=0.01, \zeta_1=0.2$)
 Fig. 3 The largest Lyapunov exponents and the corresponding isohypses

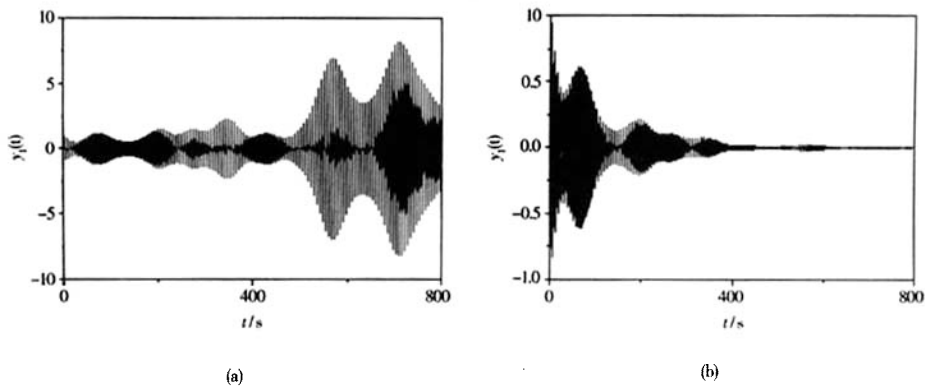


图 4 数值积分结果
 Fig. 4 The results of numerical integration

其结果见图 4(b),从图中可看出数值积分计算与最大 Lyapunov 指数计算间存在较好的一致性。

图 5 显示当增大窄带随机激励带宽 γ 时系统的最大 Lyapunov 指数及等高线图,从图中发现,图 5

(b) $\lambda=0$ 曲线较对应的图 2(b) $\lambda=0$ 曲线在底部略微平坦些,这说明在所计算的数值范围内,随机激励带宽的增大在小幅随机激励时可略微增大系统的稳定性。

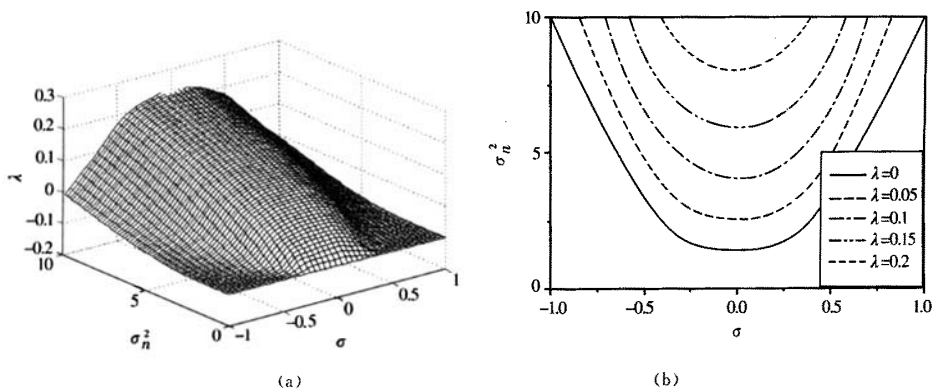


图 5 最大 Lyapunov 指数及等高线图($\gamma=0.05, \zeta_1=0.1$)
 Fig. 5 The largest Lyapunov exponents and the corresponding isohypses

4 结束语

受轴向基础窄带随机激励悬臂梁非动力学的研究目前较少,而本文在概率一意义下分析计算了系统主参激共振时平凡响应的最大 Lyapunov 指数及随机稳定性问题,使得对其复杂的非线性动力行为的观察与研究成为可能.

研究表明,增大阻尼比可有效降低系统平凡响应几乎确定不稳定区域,而增大窄带随机激励带宽只对小幅激励时的几乎确定不稳定区域略微有效.

直接数值积分验证了所获最大 Lyapunov 指数结果的正确性.

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STOCHASTIC STABILITY OF CANTILEVER BEAMS UNDERGOING AXIALLY NARROW-BAND RANDOM EXCITATION

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Abstract A dynamic modeling of an axially narrow-band random oscillating cantilever beam was presented. The motion equations for the axially oscillating cantilever beam were derived by using the Kane's equation combining with Rayleigh-Ritz method. The method of multiple scales was used to solve directly the nonlinear differential equations and to derive the nonlinear modulation equation for the principal parametric resonance. The largest Lyapunov exponent was numerically obtained to determine the almost sure stability or instability of the trivial response of the system, and the validity of the stability was verified by direct numerical integration of the equation.

Key words cantilever beam, principal parametric resonance, narrow-band random excitation, largest Lyapunov exponent, stochastic stability