

一类强非线性二阶微分方程的多模态近似解析解研究*

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摘要 利用自治力学系统的哈密顿函数为守恒量的性质,提出一种求非线性二阶微分方程多模态近似解析解的方法,称为哈密顿函数法.首先,介绍哈密顿函数法求多模态近似解的基本理论.其次,以质点在旋转的抛物线上运动为模型建立强非线性二阶微分方程.最后,用哈密顿函数法求得在给定初始条件和参数下强非线性二阶微分方程的三模态近似解析解表达式,作出三模态近似解析解的解曲线,并与直接用 Mathematica 软件作出的解曲线进行比较,讨论三模态近似解析解的精确性.结果表明:用哈密顿函数法求得的三模态近似解析解的解曲线与直接用 Mathematica 软件作出的解曲线十分吻合.

关键词 强非线性二阶微分方程, 多模态近似解析解, 哈密顿函数法

DOI: 10.6052/1672-6553-2019-067

引言

实际问题中出现的微分方程常常是非线性的,而且有许多属于强非线性的,求解非线性微分方程的近似解析解对于研究力学系统的运动特性与规律至关重要,一直受到众多学者的关注^[1-20].一些学者提出了不少有效的求解方法,如谐平衡法^[1-3],能量平衡法^[4-6],同伦渐近法^[7-9],Adomian 分解法^[10-13],哈密顿函数法^[14-17]等.在非惯性转动参照系中研究力学体系的运动时,常常会遇到一类分子分母均含非线性项的强非线性微分方程^[18],求解比较复杂或存在发散性问题,若与此类方程相应的哈密顿函数不显含时间,则属自治力学系统,其哈密顿函数是一守恒量.若运动具有周期性,则可以假设其近似解析解是多个谐振动的迭加,即多模态解,利用哈密顿函数是守恒量的性质,可以求得多模态近似解析解中各模态的振幅与基频,从而求得多模态近似解析解表达式,这种利用哈密顿函数是守恒量的性质求解的方法称哈密顿函数法,为讨论多模态近似解析解的精确性,可以在给定初始条件和参数下,在同一平面上作出强非线性二阶微分方程的多模态近似解析解的解曲线与直接用 Mathematica 软件作出的解曲线进行比较.本文用上述方法得到了强非线性二阶微分方程的三模态近

似解析解表达式,并讨论其精确性.结果表明:强非线性二阶微分方程的三模态近似解析解的解曲线与直接用 Mathematica 软件作出的解曲线十分吻合.

1 哈密顿函数法求多模态近似解析解的基本理论

若力学系统的哈密顿函数中不显含时间 t , 则其哈密顿函数 H 是一守恒量. 设某力学系统有 s 个自由度, 其哈密顿函数 H 可表示成

$$H = H(q_\alpha, \dot{q}_\alpha) \quad (\alpha = 1, 2, \dots, s) \quad (1)$$

其中, q_α, \dot{q}_α 分别表示广义坐标和广义速度. 设初始条件为

$$q_{\alpha 0} = A_\alpha, \dot{q}_{\alpha 0} = 0 \quad (\alpha = 1, 2, \dots, s) \quad (2)$$

若力学系统作周期运动, 可设其多模态近似解析解为

$$q_\alpha = \sum_{i=1}^n a_{\alpha i} \cos(i\omega t) \quad (\alpha = 1, 2, \dots, s) \quad (3)$$

则各模态的振幅满足

$$\sum_{i=1}^n a_{\alpha i} = A_\alpha \quad (\alpha = 1, 2, \dots, s) \quad (4)$$

哈密顿函数 H 在整个运动过程中均保持不变, 则其在 $0 - \frac{T}{4}$ 时间内对时间积分后仍是一不变量^[15]. 将式(3)代入式(1), 并令

2018-04-29 收到第1稿, 2019-04-09 收到修改稿.

* 国家自然科学基金资助项目(11472177)

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$$\bar{H} = \int_0^{\frac{\pi}{4}} H(q_\alpha, \dot{q}_\alpha) dt = \frac{1}{4} HT = \frac{\pi}{2\omega} H = \text{const} \quad (5)$$

则有^[15]

$$\frac{\partial}{\partial a_{\alpha i}} \left(\frac{\partial \bar{H}}{\partial (1/\omega)} \right) = 0 \quad (\alpha=1,2,\dots,s, i=1,2,\dots,n.) \quad (6)$$

联立式(4)和式(6),能解得各模态的振幅 $a_{\alpha i}$ 和振动的基频 ω , 将求得的 $a_{\alpha i}, \omega$ 代入式(3)就能得到多模态近似解析解的表达式.

2 强非线性二阶微分方程的建立

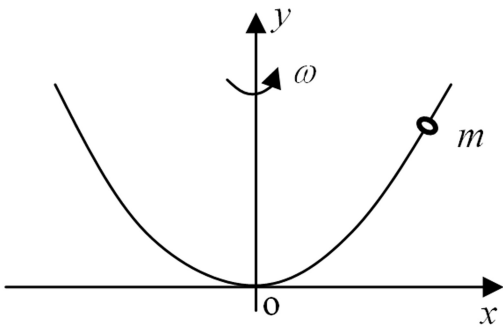


图1 匀速转动的抛物线型金属丝

Fig.1 A parabolic wire rotating at a constant speed

如图1所示,一质量为 m 的光滑小环,套在一光滑的抛物线型金属丝上,并可沿着金属丝滑动,抛物线型金属丝以角速度 ω 绕轴匀速转动.设抛物线的方程为 $x^2 = 4py$, 则小环在 x 方向的运动微分方程为^[18]

$$\ddot{x} + \frac{xx^2}{x^2 + 4p^2} + \frac{bx}{x^2 + 4p^2} = 0 \quad (7)$$

其中,

$$b = 4p^2 \left(\frac{g}{2p} - \omega^2 \right) = 4p^2 (\omega_e^2 - \omega^2) \quad (8)$$

$\omega_e = \sqrt{\frac{g}{2p}}$ 为小环保持相对平衡时金属丝的转动角速度,当 $\omega < \omega_e$ 时, $b > 0$; 当 $\omega > \omega_e$ 时, $b < 0$; 当 $\omega = \omega_e$ 时, $b = 0$.

3 强非线性二阶微分方程的三模态近似解析解及其精确性分析

方程(7)中第二项的分子分母中均含非线性项,属于强非线性二阶微分方程.用一般的方法很难求得解析近似解.文献[7]用同伦渐近法得到其近似解析解,但求解过程十分繁复,不宜推广应

用,文献[14]用哈密顿函数法只求得其单模态近似解析解,且没有对近似解进行比较研究,现用哈密顿函数方法求得其三模态近似解析解,并通过作图与数值解进行比较,分析多模态近似解的精确性.

与方程(7)相应的哈密顿函数为^[18]

$$H = \frac{1}{2} \dot{x}^2 + \frac{1}{8p^2} x^2 \dot{x}^2 + \frac{b}{8p^2} x^2 \quad (9)$$

设系统的三模态近似解为

$$x = a_1 \cos \omega t + a_2 \cos 3\omega t + a_3 \cos 5\omega t \quad (10)$$

初始条件为

$$x_0 = A, \quad \dot{x}_0 = 0 \quad (11)$$

则各模态振幅满足

$$a_1 + a_2 + a_3 = A \quad (12)$$

将式(10)代入式(9)后并按式(5)积分,有

$$\begin{aligned} \bar{H} = & \frac{1}{8} (a_1^2 + 9a_2^2 + 25a_3^2) \omega \pi + \frac{b}{32p^2} (a_1^2 + a_2^2 + a_3^2) \frac{\pi}{\omega} \\ & + \frac{1}{128p^2} (a_1^4 + 9a_2^4 + 25a_3^4 + 20a_1^2 a_2^2 + 52a_1^2 a_3^2 \\ & + 68a_2^2 a_3^2 - 8a_1^3 a_2 - 40a_1^2 a_2 a_3 - 80a_1 a_2^2 a_3) \omega \pi \end{aligned} \quad (13)$$

将式(13)代入式(6)得

$$\begin{aligned} \frac{\partial}{\partial a_1} \left(\frac{\partial \bar{H}}{\partial (1/\omega)} \right) = & - \left[\frac{1}{4} a_1 + \frac{1}{32p^2} (a_1^3 + 10a_1 a_2^2 \right. \\ & + 26a_1 a_3^2 - 6a_1^2 a_2 - 20a_1 a_2 a_3 - \\ & \left. 20a_2^2 a_3) \right] \pi \omega^2 + \frac{b}{16p^2} a_1 \pi = 0 \end{aligned} \quad (14)$$

$$\begin{aligned} \frac{\partial}{\partial a_2} \left(\frac{\partial \bar{H}}{\partial (1/\omega)} \right) = & - \left[\frac{9}{4} a_2 + \frac{1}{32p^2} (9a_2^3 + 10a_1^2 a_2 + \right. \\ & \left. 34a_2 a_3^2 - 2a_1^3 - 10a_1^2 a_3 \right. \\ & \left. - 40a_1 a_2 a_3) \right] \pi \omega^2 + \frac{b}{16p^2} a_2 \pi = 0 \end{aligned} \quad (15)$$

$$\begin{aligned} \frac{\partial}{\partial a_3} \left(\frac{\partial \bar{H}}{\partial (1/\omega)} \right) = & - \left[\frac{25}{4} a_3 + \frac{1}{32p^2} (25a_3^3 + 26a_1^2 a_3 \right. \\ & \left. + 34a_2^2 a_3 - 10a_1^2 a_2 \right. \\ & \left. - 20a_1 a_2^2) \right] \pi \omega^2 + \frac{b}{16p^2} a_3 \pi = 0 \end{aligned} \quad (16)$$

联立式(12)和式(14)~(16)四式,可以解得 a_1, a_2, a_3, ω 等4个未知量.由于式(12)和式(14)~(16)四式组成的是非线性代数方程组,比较复杂,可借助数学软件求解.讨论用哈密顿函数法得到的

近似解析解的精确性,采用的方法是将用哈密顿函数法得到的近似解析解曲线与直接用 Mathematica 软件得到的数值解曲线进行比较.为方便,现设 $A=1$, $p=b=0.25$,则式(12)和式(14)~(16)可简化为

$$a_1+a_2+a_3=1 \quad (17)$$

$$a_1-(a_1+2a_1^3+20a_1a_2^2+52a_1a_3^2-12a_1^2a_2-40a_1a_2a_3-40a_2^2a_3)\omega^2=0 \quad (18)$$

$$a_2-(9a_2+18a_2^3+20a_2^2a_1+68a_2a_3^2-4a_1^3-20a_1^2a_3-80a_1a_2a_3)\omega^2=0 \quad (19)$$

$$a_3-(25a_3+50a_3^3+52a_1^2a_3+68a_2^2a_3-20a_1^2a_2-40a_1a_2^2)\omega^2=0 \quad (20)$$

由式(17)~(20),可解得

$$a_1=-0.008246354, a_2=0.074772601, a_3=0.933473753, \omega=0.120425784. \quad (21)$$

将式(21)代入式(10),得三模态近似解析解为

$$x=-0.008246354\cos\omega t+0.074772601\cos 3\omega t+0.933473753\cos 5\omega t \quad (22)$$

其中, $\omega=0.120425784$.

根据式(22)可作出解曲线如图2中红线所示,图2中的黑线是根据式(7)直接用 Mathematica 软件作出的同参数下的解曲线.图2表明,根据哈密顿函数法得到的三模态近似解析解作出的解曲线与直接用 Mathematica 软件作出的解曲线十分吻合,特别是在第一个周期内,其振幅和周期的偏差很小,说明用哈密顿函数法得到多模态解去近似作周期运动的强非线性运动问题是一种十分有效的方法.

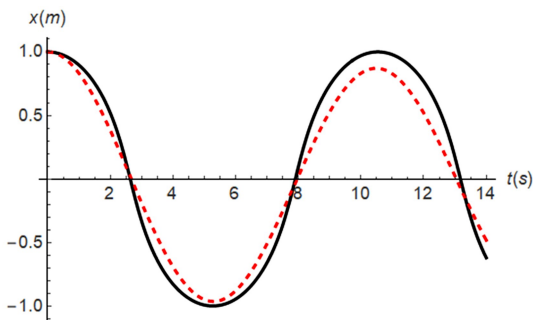


图2 哈密顿函数法得到三模态近似解曲线—红线,数值解曲线—黑线
Fig.2 The red dashed: three modes approximate solution obtained by Hamilton function method, The solid: the numerical solution

4 结论

作周期运动的力学系统可假设其解是多个不同频率谐振动的迭加,即多模态解.本文利用自治力学系统的哈密顿函数是一守恒量的性质,求得近似解析解各模态的振幅和振动的基频.在给定的初始条

件和参数下,在同一平面上作出强非线性微分方程的三模态近似解析解曲线与直接用 Mathematica 软件作出的解曲线进行比较(如图2),结果表明:三模态近似解析解的解曲线与直接用 Mathematica 软件作出的解曲线十分吻合.用哈密顿函数法求强非线性微分方程的近似解,思路简单清晰,物理意义明确,结果精确有效,可操作性强,值得推广应用.

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STUDY ON MULTIMODE APPROXIMATE ANALYTICAL SOLUTION OF A CLASS OF STRONGLY NONLINEAR SECOND ORDER DIFFERENTIAL EQUATIONS*

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Abstract Hamiltonian is a conservation quantity of autonomous mechanical system, and the multimode approximate analytical solutions of a nonlinear second order differential equation can be obtained by using characteristics of the conservation quantity, which is called Hamiltonian method. First, the basic theory of Hamiltonian method was introduced. Second, a strong nonlinear second order differential equation for the motion of a particle on a rotating parabola was established. Finally, the three-mode approximate analytical solutions of the strong nonlinear second order differential equation were obtained under given initial conditions and parameters. The approximate solution using Hamiltonian method was verified by the numerical solution using Mathematics software, which showed that the approximate solution is in good agreement with the numerical one.

Key words strongly nonlinear second order differential equation, multimode approximate analytical solutions, Hamiltonian method