运动车辆梁模型的横向振动频率及模态*

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摘要 将运动车辆的车身模型化为 Euler – Bernoulli 梁,车轮模型化为梁两端边界处的弹性不等的弹簧,形 成半车模型.通过复模态分析法研究平滑路面上移动车体的横向振动特性,给出车体横向振动的频率方程 以及模态的表达式,通过数值方法求解系统固有频率以及模态函数.并通过数值算例研究车辆运行速度、车 体刚度、轮胎弹性系数对车体横向振动的前两阶固有频率以及相应模态的影响.

关键词 半车模型, Euler - Bernoulli 梁, 横向振动, 固有频率, 模态

引 言

车辆动力学是以数学力学模型为基础并结合现代控制论,研究汽车垂直动力学、横向动力学和纵向动力学,及汽车在道路不平激励和侧风作用时的动力学性能.长期以来,吸引着众多学者关注^[1-5].

前期的车辆动力学研究均是将车体模型化为 刚体,忽略车体本身的弹性.研究方法主要是将汽 车的四分之一模型化为质量 - 弹簧阻尼 - 质量 -弹簧系统的两自由度问题^[1,2],即四分之一汽车模 型.将汽车模型化为车身内部含有两个弹性支撑的 刚性质量块的二维模型^[3,4],即半车模型.还有将 汽车模型化为四角支撑边界的刚性板的三维模 型^[5],即整车模型.为了研究考虑车体弹性的振动 问题,本文将车体模型化为轴向运动的 Euler - Bernoulli 梁,将车轮对车体的作用模型化为梁边界处 的弹性支撑,从而形成半车模型,并研究在平滑路 面上,即无外激励时的振动特性.

对于轴向运动的 Euler – Bernoulli 梁的研究有 着丰富的成果. Simpson 研究了固定边界条件下,轴 向运动梁的固有频率及模态^[6]. Öz 等给出了两端简 支^[7]的轴向运动梁的频率和模态. 王洪霞等研究了 含表面裂纹简支梁的非线性振动频率^[8]. 马连生等 研究比较了简支边界下 Euler – Bernoulli 梁和其他 梁模型的不同^[9]. 李彪等给出了混杂边界的轴向运 动梁的横向振动频率和模态^[10]. 相关于垂直方向弹 性支撑边界的轴向运动系统的研究还未见报道.本 文研究以车轮为工程背景的弹性支撑边界轴向运动 的 Euler – Bernoulli 梁的横向振动频率以及模态.

1 控制方程

水平路面移动车辆的半车模型如图 1 所示. 其 中,设其水平平衡位置为 X 轴, T 为时间坐标. 将汽 车车身简化为轴向运动的 Euler – Bernoulli 梁,即 以速度 V 沿着 X 方向作匀速运动,梁的长度为 L, 密度为ρ,等截面的横截面积为 A,弹性模量为 E, 惯性矩为 I,并假设两端有初始拉力 P,以弹簧为模 型的车轮的弹性系数分别为 K₁ 和 K₂.





根据达朗伯原理,分析梁微元横向受力,得到 车身横向振动的无量纲控制方程

$$u_{,u} + 2Vu_{,xt} + V^{2}u_{,xx} - u_{,xx} + k_{f}^{2}u_{,xxx} = 0$$
 (1)
梁两端由弹簧支撑的边界条件为
$$u_{,xx}(0,t) = 0, u_{,xxx}(0,t) - k_{1}u(0,t) = 0$$

$$u_{,xx}(1,t) = 0, u_{,xxx}(1,t) - k_{2}u(1,t) = 0$$

其中无量纲变量及参数定义加下

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$$u = \frac{U}{L}, t = T_{\sqrt{\frac{P}{\rho A L^2}}}, v = V_{\sqrt{\frac{\rho A}{P}}}, x = \frac{X}{L},$$
$$k_f^2 = \frac{EI}{PL^2}, k_1 = \frac{K_1 L}{P}, k_2 = \frac{K_2 L}{P}$$
(3)

式中,下角标表示对下角标变量的求导.

2 固有频率及模态函数

在文献[7,10]中,对于不同边界下轴向运动 梁的固有频率以及模态求解,都是通过复模态的形 式引入并加以处理的.受此启发,这里也假设梁的 模态函数表现为复数形式,即假设梁单元横向振动 位移满足

$$u(x,t) = \phi_n(x)e^{i\omega_n t} + \phi_n(x)e^{-i\omega_n t}$$
(4)

其中,ω_n,φ_n分别表示轴向运动梁的第 n 固有频率 和模态函数.把解的形式代入(1)式中,得到

 $-\omega_n^2 \phi_n + 2iv\omega_n \phi'_n + (v^2 - 1)\phi''_n + k_f^2 \phi_n^{(4)} = 0$ (5) 由于频率方程(5)是四阶常微分方程,可设方程的 本征方程的根为 $i\beta_{j_n}(j = 1, 2, 3, 4)$,即 β_{j_n} 满足实系 数4次代数本征方程

$$k_r^2 \beta_{jn}^4 - (v^2 - 1)\beta_{jn}^2 - 2v\omega_n \beta_{jn} - \omega_n^2 = 0$$
(6)
则轴向运动梁的模态函数可写做

 $\phi_n(x) = C_{1n}(e^{i\beta_{1n}x} + C_{2n}e^{i\beta_{2n}x} + C_{3n}e^{i\beta_{3n}x} + C_{4n}e^{i\beta_{4n}x}) \quad (7)$ 对模态函数归—化^[7]

$$\phi_n(x) = \frac{1}{\sqrt{\rho}} \left(e^{i\beta_{1n^x}} + C_{2n} e^{i\beta_{2n^x}} + C_{3n} e^{i\beta_{3n^x}} + C_{4n} e^{i\beta_{4n^x}} \right) \quad (8)$$

将(4)式代入边界条件(2)式中,化简得到

$$\phi_n^{''}(0) = 0, \phi_n^{'''}(0) - k_1 \phi_n(0) = 0,$$

$$\phi_n^{''}(1) = 0, \phi_n^{'''}(1) - k_2 \phi_n(1) = 0$$
(9)

再将(8)式代入(9)式中,得到线性方程组

$$-\left(\beta_{1n}^{2}+C_{2n}\beta_{2n}^{2}+C_{3n}\beta_{3n}^{2}+C_{4n}\beta_{4n}^{2}\right)=0$$
(10)
$$-\left(i\beta_{1n}^{3}+k_{1}\right)-C_{2n}\left(i\beta_{2n}^{3}+k_{1}\right)-C_{3n}\left(i\beta_{3n}^{3}+k_{1}\right)-C_{4n}\left(i\beta_{4n}^{3}+k_{1}\right)=0$$
(11)

$$\beta_{1n}^2 e^{i\beta_{1n}} + C_{2n}\beta_{2n}^2 e^{i\beta_{2n}} + C_{3n}\beta_{3n}^2 e^{i\beta_{3n}} + C_{4n}\beta_{4n}^2 e^{i\beta_{4n}} = 0$$

$$-C_{2n}(i\beta_{2n}^{3}+k_{2})e^{i\beta_{2n}}-C_{3n}(i\beta_{3n}^{3}+k_{2})e^{i\beta_{3n}}-C_{4n}(i\beta_{4n}^{3}+k_{2})e^{i\beta_{4n}}-(i\beta_{1n}^{3}+k_{2})e^{i\beta_{1n}}=0$$
(13)
联合以上4个线性方程,写成矩阵

$$\begin{pmatrix} \beta_{1n}^{2} & \beta_{2n}^{2} & \beta_{3n}^{2} & \beta_{4n}^{2} \\ (i\beta_{1n}^{3} + k_{1}) & (i\beta_{2n}^{3} + k_{1}) & (i\beta_{3n}^{3} + k_{1}) & (i\beta_{4n}^{3} + k_{1}) \\ \beta_{1n}^{2}e^{i\beta_{1n}} & \beta_{2n}^{2}e^{i\beta_{2n}} & \beta_{3n}^{2}e^{i\beta_{3n}} & \beta_{4n}^{2}e^{i\beta_{4n}} \\ (i\beta_{1n}^{3} + k_{2})e^{i\beta_{1n}} & (i\beta_{2n}^{3} + k_{2})e^{i\beta_{2n}} & (i\beta_{3n}^{3} + k_{2})e^{i\beta_{3n}} & (i\beta_{4n}^{3} + k_{1})e^{i\beta_{4n}} \end{pmatrix} \begin{pmatrix} 1 \\ C_{2n} \\ C_{3n} \\ C_{4n} \end{pmatrix} (14)$$

数表达式为

 $C_{4n} = -\frac{\beta_{1n}^2 + C_{2n}\beta_{2n}^2 + C_{3n}\beta_{3n}^2}{\beta_{4n}^2}$

根据(14)式有非零解的充要条件,即系数行列式 为零,可以解出 β_{jn} ,j=1,2,3,4.结合频率方程(6) 式即可解得固有频率及模态函数,其中模态函数系 C_{2n} =

$$\frac{-i\beta_{1n}^{2}\beta_{3n}^{2}\beta_{4n}^{2}\left[\left(e^{i\beta_{3n}}-e^{i\beta_{4n}}\right)\beta_{1n}+\left(-e^{i\beta_{1n}}+e^{i\beta_{4n}}\right)\beta_{3n}+\left(e^{i\beta_{1n}}-e^{i\beta_{3n}}\right)\beta_{4n}\right]+k_{1}\left\{-\left(e^{i\beta_{3n}}-e^{i\beta_{4n}}\right)\beta_{3n}^{2}\beta_{4n}^{2}+\beta_{1n}^{2}\left[\left(-e^{i\beta_{1n}}+e^{i\beta_{3n}}\right)\beta_{3n}^{2}+\left(e^{i\beta_{1n}}-e^{i\beta_{4n}}\right)\beta_{4n}^{2}\right]\right\}}{i\beta_{2n}^{2}\beta_{3n}^{2}\beta_{4n}^{2}\left[\left(\left(e^{i\beta_{3n}}-e^{i\beta_{4n}}\right)\beta_{2n}\right)+\left(-e^{i\beta_{2n}}+e^{i\beta_{4n}}\right)\beta_{3n}+\left(e^{i\beta_{2n}}-e^{i\beta_{3n}}\right)\beta_{4n}\right]+k_{1}\left\{\left(e^{i\beta_{3n}}-e^{i\beta_{4n}}\right)\beta_{3n}^{2}\beta_{4n}^{2}+\beta_{2n}^{2}\left[\left(e^{i\beta_{2n}}-e^{i\beta_{3n}}\right)\beta_{3n}^{2}-\left(e^{i\beta_{2n}}-e^{i\beta_{4n}}\right)\beta_{4n}^{2}\right]\right\}}$$

$$(16)$$

$$C_{3n} = \frac{-i\beta_{1n}^{2}\beta_{2n}^{2}\beta_{4n}^{2}\left[(e^{i\beta_{2n}}-e^{i\beta_{4n}})\beta_{1n} + (-e^{i\beta_{1n}}+e^{i\beta_{4n}})\beta_{2n} + (e^{i\beta_{1n}}-e^{i\beta_{2n}})\beta_{4n}\right] + k_{1}\left\{-(e^{i\beta_{2n}}-e^{i\beta_{4n}})\beta_{2n}^{2}\beta_{4n}^{2} + \beta_{1n}^{2}\left[(-e^{i\beta_{1n}}+e^{i\beta_{2n}})\beta_{2n}^{2} + (e^{i\beta_{1n}}-e^{i\beta_{4n}})\beta_{4n}^{2}\right]\right\}} \\ -i\beta_{2n}^{2}\beta_{3n}^{2}\beta_{4n}^{2}\left[((e^{i\beta_{3n}}-e^{i\beta_{4n}})\beta_{2n}) + (-e^{i\beta_{2n}}+e^{i\beta_{4n}})\beta_{3n} + (e^{i\beta_{2n}}-e^{i\beta_{3n}})\beta_{4n}\right] + k_{1}\left\{-(e^{i\beta_{3n}}-e^{i\beta_{4n}})\beta_{3n}^{2}\beta_{4n}^{2} + \beta_{2n}^{2}\left[(-e^{i\beta_{2n}}+e^{i\beta_{3n}})\beta_{3n}^{2} - (e^{i\beta_{2n}}-e^{i\beta_{4n}})\beta_{4n}^{2}\right]\right\}$$

展开方程(14)式的系数行列式,得到频率方程

$$e^{i\beta_{4n}}[e^{i\beta_{3n}}(k_1\beta_{2n}^2-k_1\beta_{1n}^2+i\beta_{1n}^3\beta_{2n}^2-i\beta_{1n}^2\beta_{2n}^3)(k_2+i\beta_{3n}^3) -$$

 $e^{i\beta_{2n}}(k_2+i\beta_{2n}^3)(k_1\beta_{3n}^2-k_1\beta_{1n}^2+i\beta_{1n}^3\beta_{3n}^2-i\beta_{1n}^2\beta_{3n}^3) +$
 $e^{i\beta_{1n}}(k_2+i\beta_{1n}^3)(k_1\beta_{3n}^2-k_1\beta_{2n}^2+i\beta_{2n}^3\beta_{3n}^2 -$
 $i\beta_{2n}^2\beta_{3n}^3)]\beta_{4n}^2 - e^{i\beta_{3n}}\beta_{3n}^2[e^{i\beta_{4n}}(k_2+i\beta_{4n}^3)(k_1\beta_{2n}^2 -$
 $k_1\beta_{1n}^2+i\beta_{1n}^3\beta_{2n}^2-i\beta_{1n}^2\beta_{2n}^3) - e^{i\beta_{2n}}(k_2+i\beta_{2n}^3)(k_1\beta_{4n}^2 -$

$$\begin{split} & k_{1}\beta_{1n}^{2} + i\beta_{1n}^{3}\beta_{3n}^{2} - i\beta_{1n}^{2}\beta_{4n}^{3} \right) + e^{i\beta_{1n}}(k_{2} + i\beta_{1n}^{3})(k_{1}\beta_{4n}^{2} - k_{1}\beta_{2n}^{2}\beta_{4n}^{2} - i\beta_{2n}^{2}\beta_{4n}^{3})] + e^{i\beta_{2n}}\beta_{2n}^{2}[e^{i\beta_{4n}}(k_{2} + i\beta_{3n}^{3})(k_{1}\beta_{3n}^{2} - k_{1}\beta_{1n}^{2}\beta_{3n}^{2})] + e^{i\beta_{2n}}\beta_{3n}^{2}(e^{i\beta_{4n}}(k_{2} + i\beta_{3n}^{3}))(k_{1}\beta_{4n}^{2} - k_{1}\beta_{1n}^{2}\beta_{3n}^{2} - i\beta_{1n}^{2}\beta_{3n}^{3})) - e^{i\beta_{3n}}(k_{2} + i\beta_{3n}^{3})(k_{1}\beta_{4n}^{2} - k_{1}\beta_{1n}^{2} + i\beta_{3n}^{3}\beta_{4n}^{2} - i\beta_{1n}^{2}\beta_{4n}^{3})) + e^{i\beta_{1n}}(k_{2} + i\beta_{1n}^{3})(k_{1}\beta_{4n}^{2} - k_{1}\beta_{3n}^{2} + i\beta_{3n}^{3}\beta_{4n}^{2} - i\beta_{2n}^{2}\beta_{4n}^{3})] - e^{i\beta_{1n}}\beta_{1n}^{2}\left[e^{i\beta_{4n}}(k_{1}\beta_{3n}^{2} - k_{1}\beta_{2n}^{2} + i\beta_{3n}^{2}\beta_{3n}^{2} - i\beta_{2n}^{2}\beta_{3n}^{3})(k_{2} + i\beta_{3n}^{2}\beta_{4n}^{2} - i\beta_{3n}^{2}\beta_{3n}^{2} - i\beta_{3n}^{2}\beta_{3n}^{2})(k_{2} + i\beta_{3n}^{2}\beta_{3n}^{2} - i\beta_{2n}^{2}\beta_{3n}^{3})(k_{2} + i\beta_{3n}^{2}\beta_{3n}^{2})(k_{2} + i\beta_{3n}^{2})(k_{2} + i\beta_{3n}^{2}$$

(12)

(15)

$$i\beta_{4n}^{3} - (k_{2} + i\beta_{3n}^{3})e^{i\beta_{3n}}(k_{1}\beta_{4n}^{2} - k_{1}\beta_{2n}^{2} + i\beta_{2n}^{3}\beta_{3n}^{2} - i\beta_{2n}^{2}\beta_{4n}^{3}) + e^{i\beta_{2n}}(k_{2} + i\beta_{2n}^{3})(k_{1}\beta_{4n}^{2} - k_{1}\beta_{3n}^{2} + i\beta_{3n}^{3}\beta_{4n}^{2} - i\beta_{3n}^{2}\beta_{4n}^{3})] = 0$$
(18)

3 数值结果

计算中取车身的弹性模量为2.1×10¹¹ Pa,密度 为7850 kg/m³,配合系统的物理参数可得到系统的 各个无量纲量.图2给出了对应于不同的车身弯曲 刚度时,车体系统的前两阶固有频率.图2中车轮弹 性系数 k₁=1000,k₂=1000.观察发现,系统前两阶的 固有频率值均随车身刚度的增大而增大,比较图(a) 和(b)可以发现,在一定速度范围内,车身前两阶横 向振动固有频率随车体移动速度的增大而减小.



图 2 固有频率随速度和车身刚度的变化

Fig. 2 Natural frequencies changing with axial speed and



图 3 速度和车轮弹性系数对固有频率的影响 Fig. 3 Natural frequencies changing with axial speed and stiffness of the tires

图 3 给出了两个车轮的弹性系数 k₁ 和 k₂ 取相 等的不同值时,车身前两阶横向振动固有频率. 图 3 中车身的弯曲刚度 k_f = 0.8.观察图 3 发现,车身 的前两阶横向振动固有频率的值随任意一个车轮 的弹性系数增大而增大,而且频率增大的幅度随弹 性系数的增大而减小.同时,比较还发现,第二阶固 有频率受到两个车轮弹性系数的影响要大于第一 阶固有频率.同时,随着车体移动速度的增大,两个 车轮弹性系数对车身前两阶横向振动固有频率的 影响将变小.



Fig. 4 Natural frequencies changing with stiffness of the tire

图 4 显示了固定一个车轮弹性系数 $k_1 = 1000$ 不变,而另一个系数 k_2 取不同值时,车身前两阶横 向振动固有频率.图 4 中计算参数取值为车身弯曲 刚度 $k_f = 0.8$,车体移动速度 v = 1.观察图 4 所呈现 的现象表明,如果固定其中一个车轮的弹性系数 k_1 不变,当车轮弹性系数 k_2 不大时,车身前两阶横向 振动固有频率的值随弹性系数 k_2 的增大而增大, 同时发现,随 k_2 的增大, k_2 对前两阶横向振动固有 频率的影响将变弱,这些现象与从图 3 中观察到的 完全一致.

图 5 和 6 分别给出了对应于前两阶固有频率 的静态和动态车身前两阶振动模态. 图 5 和 6 中车 轮弹性系数 k₁ = 1000,k₂ = 1000,k_f = 0.8,图 5 和 6 中的车体移动速度分别为 v = 0 和 v = 2. 图中实线 为实部,点划线为虚部. 观察发现,相比较于静态车 体,动态车体模态函数的实部在形状以及幅度上都 发生了明显的变化,而虚部的幅度发生了变化,但 是形状变化不明显,即由于速度的影响,使得模态 1.0



图5 静态前两阶模态

Fig. 5 The first two modes for static state



图 6 动态前两阶模态

Fig. 6 The first two modes for dynamic



图 7 较小车轮弹性系数的动态前两阶模态 Fig. 7 The first two modes for small stiffness of the tire

图 7 中车轮弹性系数 k₁ = 100,k₂ = 100,k_f = 0. 8,v = 2. 观察发现,相比较于图 6,较小的车轮弹性 系数的动态车体模态函数的对称性明显较差,可见 增加车轮弹性系数可以使得模态函数的对称性明 显增强.

4 结论

建立以 Euler - Bernoulli 梁为车辆车体模型、 两个弹性系数不等的弹簧为车轮模型,并以边界处 的作用引入车轮对车身作用的半车模型.考虑平滑 路面上的匀速移动车辆,即无外激励的作用.通过 传统的半解析半数值的方法,即运用复模态解析方 法分析,建立车体横向振动的本征方程、频率方程 以及模态函数的表达式,再通过数值途径计算超越 方程,得到车体横向振动的固有频率以及模特.数 值结果表明,车体的前两阶横向振动固有频率随着 车体移动速度的增大而减小,随着车体弯曲刚度以 及车轮弹性系数的增大而增大,而且随着车轮弹性 系数的增大,车体前两阶固有频率增大的幅度将逐 渐变小.通过对静态和动态模态函数的比较,发现 车体的移动速度和车轮弹性系数均对模态函数有 着明显的影响.

参考文献

- 钱长照.车桥耦合振动的摄动分析方法.动力学与控制 学报,2009,7(4):375~378 (Qian C Z. Perturbation method for coupling vibration of vehicle-bridge system. *Journal of Dynamics and Control*, 2009,7(4):375~378 (in Chinese))
- 2 Yang S P, Pan C Z. A hysteresis model for magneto-rheological damper. International Journal of Nonlinear Sciences and Numerical Simulation, 2005,6(2): 139 ~ 144
- 3 Law S S, Bu J Q, Zhu XQ, Chan SL. Vehicle axle loads identification using finite element method . *Engineering Structures*, 2004, 26: 1143 ~ 1153
- 4 吴龙,闻霞. 6 自由度半车悬架解耦及其分层振动控制的研究.汽车工程,2010,32(2):148~154 (Wu L, Wen X. A research on the decoup ling of a 6 DOF half vehicle suspension and its hierarchical vibration control. *Automotive Engineering*, 2010,32(2):148~154 (in Chinese))
- 5 ElBeheiry E M. Two tracking control problems applied to damping nonlinear vehicle vibrations. *Journal of Sound and Vibration*, 2009,321: 471 ~ 491
- 6 Simpson A. Transverse modes and frequencies of beams translating between fixed end supports. *Journal of Mechani*cal Engineering Science, 1973, 15(3): 159 ~ 164
- 7 z H R. On the vibrations of an axially traveling beam on fixed supports with variable velocity. *Journal of Sound Vibration*, 2001,239,556 ~ 564
- 8 王洪霞,李学平. 含表面裂纹简支梁的非线性振动分析. 动力学与控制学报,2010,18(2):177~181(Wang H X,Li X P. Nonlinear vibration analysis of simply supported beams with an edge crack . *Journal of Dynamics and Control*, 2010,18(2):177~181 (in Chinese))
- 9 马连生,欧志英,黄达文.不同梁理论之间简支梁特征值的 解析关系.工程力学,2006,23(10):91~95(Ma L S,Ou Z Y,Huang D W. Analytical relationships of simply-supported beam's eigenvalues using different beam theories. *Engineering Mechanics*, 2006,23(10):91~95 (in Chinese))

 李彪,丁虎,陈立群.非对称混杂边界轴向运动 Timoshenko 梁横向振动分析.固体力学学报,2009,30(6):
 565~570 (Li B, Ding H, Chen L Q. Transverse vibration of axially moving timoshenko beams with unsymmetrical hybrid supports. *Acta Mechanica Solida Sinica*, 2009, 30 (6):565 ~ 570 (in Chinese))

NATURAL FREQUENCIES AND MODEL FUNCTION OF TRANSVERSE VIBRATION OF BEAM MODAL FOR MOVING VEHICLES*

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Abstract The car's body was modeled by the Euler-Bernoulli beam model, and the car's tires were modeled by the two springs with different spring index at both ends for the half-vehicle modal. The first two natural frequencies and model function of transverse vibration of moving vehicles over a smooth road were numerically solved. The transcendental equation and the eigenfunction were established by boundary condition based on the complex model function. The first two natural frequencies and the corresponding model function were obtained numerically, and the contributions of the speed of vehicle, the flexural stiffness of the car's body and the stiffness of the tires were discussed via numerical examples.

Key words half-vehicle modal, Euler-Bernoulli beam, transverse vibration, natural frequency, modal function

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