

基于对偶变量变分原理和两端动量独立变量的保辛方法*

高强 谭述君 张洪武 钟万勰

(大连理工大学工程力学系,工业装备结构分析国家重点实验室,大连 116023)

摘要 将广义位移和动量同时用拉格朗日多项式近似,并选择积分区间两端动量为独立变量,然后基于对偶变量变分原理导出了哈密顿系统的离散正则变换和对应的数值积分保辛算法.当位移和动量的拉格朗日多项式近似阶数满足一定条件时,可以自然导出保辛算法的不动点格式.通过数值算例讨论了位移和动量采用不同阶次插值所需最少 Gauss 积分点个数,并讨论了位移插值阶数、动量插值阶数以及 Gauss 积分点个数对保辛算法精度的影响,说明了上述不动点格式恰好是一种最优格式.

关键词 变分原理, 保辛方法, 哈密顿系统, 对偶

引言

保守的动力系统可以用哈密顿(Hamilton)体系描述.哈密顿动力系统的一个主要性质是其相流保持辛结构^[1,2].一般的非线性哈密顿系统的闭合解通常难以得到,因此,数值积分方法常常是必须的选择.由于哈密顿动力系统相流保辛的特点,其近似积分的数值方法也应该设法保持相流的辛结构^[3].非线性哈密顿系统的保辛算法已有许多研究,多是基于差分近似的 Runge - Kutta 方法,见有代表性的著作^[3-7].保辛 Runge - Kutta 法的困难之处在于构造高精度算法.本文基于分析结构力学^[8-10]并利用有限元的概念,将广义位移和动量同时用拉格朗日多项式近似,并选择积分区间两端动量为独立变量,然后基于对偶变量变分原理导出了哈密顿系统的离散正则变换和对应的数值积分保辛算法.通过数值算例讨论了当位移和动量采用不同阶次多项式插值时所需最少的 Gauss 积分点个数,并讨论了位移插值阶数、动量插值阶数以及 Gauss 积分点个数对保辛算法精度的影响.

1 哈密顿系统的基本方程和变分原理

假设哈密顿动力系统有 d 个广义坐标 $\mathbf{q} = \{q_1, q_2, \dots, q_d\}^T$,其拉格朗日函数为 $L(\mathbf{q}, \dot{\mathbf{q}})$,一点表示对时间的导数,则作用量 S 为拉格朗日函数在时间域上的积分,即

$$S = \int_0^\eta L(\mathbf{q}, \dot{\mathbf{q}}) dt \quad (1)$$

哈密顿变分原理为

$$\delta S = \int_0^\eta \left(\frac{\partial L}{\partial \mathbf{q}} \delta \mathbf{q} + \frac{\partial L}{\partial \dot{\mathbf{q}}} \delta \dot{\mathbf{q}} \right) dt = \int_0^\eta \left[\frac{\partial L}{\partial \mathbf{q}} - \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\mathbf{q}}} \right) \right] \delta \mathbf{q} dt + \left. \left(\frac{\partial L}{\partial \dot{\mathbf{q}}} \right) \delta \mathbf{q} \right|_0^\eta = 0 \quad (2)$$

若认为两端的广义位移给定,即 $\delta \mathbf{q}$ 在 $t=0$ 和 $t=\eta$ 两端为零,则得到欧拉-拉格朗日方程

$$\frac{\partial L}{\partial \mathbf{q}} - \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\mathbf{q}}} \right) = 0 \quad (3)$$

欧拉-拉格朗日方程是以广义位移一类变量表示的运动方程.

对拉格朗日函数中的变量 $\dot{\mathbf{q}}$ 作勒让德(Legendre)变换,即

$$\mathbf{p} = \frac{\partial L}{\partial \dot{\mathbf{q}}} \quad (4)$$

并定义哈密顿函数

$$H(\mathbf{q}, \mathbf{p}) = \mathbf{p}^T \dot{\mathbf{q}} - L(\mathbf{q}, \dot{\mathbf{q}}) \quad (5)$$

则用对偶变量表示的作用量为

$$S = \int_0^\eta [\mathbf{p}^T \dot{\mathbf{q}} - H(\mathbf{q}, \mathbf{p})] dt \quad (6)$$

对偶变分原理为

$$\delta S = \int_0^\eta \left[(\delta \dot{\mathbf{q}})^T \mathbf{p} - (\delta \mathbf{q})^T \frac{\partial H}{\partial \mathbf{q}} \right] dt + \int_0^\eta (\delta \mathbf{p})^T \left(\dot{\mathbf{q}} - \frac{\partial H}{\partial \mathbf{p}} \right) dt = \int_0^\eta (\delta \mathbf{q})^T \left[-\dot{\mathbf{p}} - \frac{\partial H}{\partial \mathbf{q}} \right] dt + \int_0^\eta (\delta \mathbf{q})^T T \left(\dot{\mathbf{q}} - \frac{\partial H}{\partial \mathbf{p}} \right) dt + \mathbf{p}^T \delta \mathbf{q} \Big|_0^\eta = 0 \quad (7)$$

2008-12-10 收到第1稿,2008-12-30 收到修改稿.

* 国家自然科学基金(10632030,10721062,2005CB321704),辽宁省博士启动基金(20081091) 大连理工大学青年教师培养基金,大连理工大学理学基金(SFDUT07002)资助项目资助

若认为两端的广义位移给定,即 δq 在两端为零,则得到哈密顿正则方程

$$\dot{q} = \frac{\partial H}{\partial p}, \dot{p} = -\frac{\partial H}{\partial q} \quad (8)$$

哈密顿正则方程是以广义位移和动量两类对偶变量表示的运动方程. 方程(8)也可以写成如下的形式

$$\dot{v} = J \frac{\partial H}{\partial v}, v = \{q^T, p^T\}^T \quad (9)$$

其中 J 是单位辛矩阵,即

$$J = \begin{bmatrix} 0_d & I_d \\ -I_d & 0_d \end{bmatrix} \quad (10)$$

其中 0_d 和 I_d 分别是 d 维的零矩阵和单位矩阵. 哈密顿动力系统的最主要的结构特性是相流保辛.

2 基于对偶变量变分原理的近似作用量

对于有些系统,直接给出的是哈密顿函数,若要利用作用量构造保辛算法,需要首先通过勒让德变换将哈密顿函数转换为 Lagrange 函数. 但对于非线性系统,这个转化常常是不可行的,因此下面我们构造基于 Hamilton 函数的保辛算法.

假设积分区间为 $[0, \eta]$, 将积分区间等分为 m 份, 长度为 η/m , 在此时间区间内将广义坐标 $q_i(t)$ 用通过上述等分点的 m 阶 Lagrange 多项式近似. 同样, 将积分区间 $[0, \eta]$ 等分为 n 份, 长度为 η/n , 在此时间区间内将广义动量 $p_i(t)$ 用通过上述等分点的 n 阶 Lagrange 多项式近似, 即

$$q_i(t) = \sum_{k=0}^m M_k(t) q_k^i, i=1, 2, \dots, d \quad (11)$$

$$p_i(t) = \sum_{k=0}^n N_k(t) p_k^i, i=1, 2, \dots, d \quad (12)$$

其中

$$q_k^i = q^i\left(\frac{k}{m}\eta\right), p_k^i = p^i\left(\frac{k}{n}\eta\right) \quad (13)$$

$$M_k(t) = \prod_{i=0, i \neq k}^m \frac{t - i\eta/m}{\eta/m(k-i)}, k=0, 1, \dots, m$$

$$N_k(t) = \prod_{i=0, i \neq k}^n \frac{t - i\eta/n}{\eta/n(k-i)}, k=0, 1, \dots, n \quad (14)$$

则广义位移向量和动量向量可表示为

$$q(t) = \sum_{k=0}^m M_k(t) q_k, p(t) = \sum_{k=0}^n N_k(t) p_k \quad (15)$$

其中

$$q_k = \{q_k^1, q_k^2, \dots, q_k^d\}^T, k=0, 1, \dots, m$$

$$p_k = \{p_k^1, p_k^2, \dots, p_k^d\}^T, k=0, 1, \dots, n \quad (16)$$

则通过方程(6)可计算近似作用量为

$$S(q_0, q_1, \dots, q_m; p_0, p_1, \dots, p_n) = \int_0^\eta [p^T \dot{q} - H(q, p)] dt \quad (17)$$

下面考虑基于近似作用量 $S(q_0, q_1, \dots, q_m; p_0, p_1, \dots, p_n)$ 的正则变换和保辛算法.

3 以 p_0 和 p_n 为独立变量的保辛算法

根据方程(7), 如果 q 和 p 在域内满足哈密顿正则方程, 则作用量的微商为

$$p_n^T dq_m - p_0^T dq_0 = dS \quad (18)$$

如果选择 p_0 和 p_n 为独立变量, 根据方程(18)有

$$-q_m^T dp_n + q_0^T dp_0 = d(p_0^T q_0 - p_n^T q_m + S) \quad (19)$$

根据方程(17), S 是 q_0, q_1, \dots, q_m 和 p_0, p_1, \dots, p_n 的函数, 因此有

$$d(p_0^T q_0 - p_n^T q_m + S) = \frac{\partial S}{\partial q_1} dq_1 + \dots + \frac{\partial S}{\partial q_{m-1}} dq_{m-1} + \frac{\partial S}{\partial p_1} dp_1 + \dots + \frac{\partial S}{\partial p_{n-1}} dp_{n-1} + \left(\frac{\partial S}{\partial p_n} - q_m\right) dp_n + \left(\frac{\partial S}{\partial p_0} + q_0\right) dp_0 + \left(p_0 + \frac{\partial S}{\partial q_0}\right) dq_0 + \left(\frac{\partial S}{\partial q_m} - p_n\right) dq_m \quad (20)$$

由于选则 p_0 和 p_n 为独立变量, 那么正则变换要求 $d(p_0^T q_0 - p_n^T q_m + S)$ 为关于独立变量 p_0 和 p_n 的全微商, 即要求

$$\frac{\partial S}{\partial q_0} + p_0 = 0, \frac{\partial S}{\partial q_m} - p_n = 0, \frac{\partial S}{\partial q_i} = 0, i=1, 2, \dots, m-1,$$

$$\frac{\partial S}{\partial p_i} = 0, i=1, 2, \dots, n-1 \quad (21)$$

上式为关于 $(m+n+2)d$ 个变量 $q_0, q_1, \dots, q_m; p_0, p_1, \dots, p_n$ 的方程, 而方程的个数为 $(m+n)d$ 个, 因此可解出

$$q_0 = \bar{q}_0(p_0, p_n), q_1 = \bar{q}_1(p_0, p_n), \dots, q_m = \bar{q}_m(p_0, p_n)$$

$$p_0 = \bar{p}_0(p_0, p_n), p_1 = \bar{p}_1(p_0, p_n), \dots,$$

$$p_{n-1} = \bar{p}_{n-1}(p_0, p_n) \quad (22)$$

将方程(22)代入 $p_0^T q_0 - p_n^T q_m + S$, 则作用量成为 p_0 和 p_n 函数, 即

$$\bar{S}(p_0, p_n) = p_0^T \bar{q}_0(p_0, p_n) - p_n^T \bar{q}_m(p_0, p_n) + S(\bar{q}_0(p_0, p_n), \bar{q}_1(p_0, p_n), \dots, \bar{q}_m(p_0, p_n); p_0, \bar{p}_1(p_0, p_n), \dots, \bar{p}_{n-1}(p_0, p_n), p_n) \quad (23)$$

那么 $\bar{S}(p_0, p_n)$ 对应的正则变换为

$$q_0 = \frac{\partial \bar{S}}{\partial p_0}, q_m = -\frac{\partial \bar{S}}{\partial p_n} \quad (24)$$

方程(24)给出了离散哈密顿正则方程. 以上是理论分析,但对于非线性问题,方程(21)不可能解析求解出(22),然后代入 $\mathbf{p}_0^T \mathbf{q}_0 - \mathbf{p}_n^T \mathbf{q}_m + S$,因此需要考虑数值求解过程.

如果将 $\mathbf{p}_0 \mathbf{q}_0 - \mathbf{p}_n \mathbf{q}_m + S$ 中的 $\mathbf{q}_0, \mathbf{q}_1, \dots, \mathbf{q}_m$ 分别用新变量 $\bar{\mathbf{q}}_0, \bar{\mathbf{q}}_1, \dots, \bar{\mathbf{q}}_m$ 代替, $\mathbf{p}_1, \mathbf{p}_2, \dots, \mathbf{p}_{n-1}$ 用分别用新变量 $\bar{\mathbf{p}}_1, \bar{\mathbf{p}}_2, \dots, \bar{\mathbf{p}}_{n-1}$ 代替,即将 S 表示为 $S(\mathbf{q}_0, \bar{\mathbf{q}}_1, \dots, \bar{\mathbf{q}}_m; \bar{\mathbf{p}}_0, \dots, \bar{\mathbf{p}}_{n-1}, \mathbf{p}_n)$,那么求解如下方程

$$\frac{\partial S}{\partial \mathbf{q}_0} + \mathbf{p}_0 = 0, \frac{\partial S}{\partial \bar{\mathbf{q}}_i} = 0, \frac{\partial S}{\partial \mathbf{q}_m} - \mathbf{p}_n = 0, \quad i = 1, 2, \dots, m-1 \quad (25)$$

$$\frac{\partial S}{\partial \mathbf{p}_0} + \bar{\mathbf{q}}_0 - \mathbf{q}_0 = 0, \frac{\partial S}{\partial \bar{\mathbf{p}}_i} = 0, \mathbf{q}_m = \bar{\mathbf{q}}_m - \frac{\partial S}{\partial \mathbf{p}_n}, \quad i = 1, 2, \dots, n-1 \quad (26)$$

与方程(24)给出的正则变换等价. 若记

$$\bar{\mathbf{q}} = \{\bar{\mathbf{q}}_0^T, \bar{\mathbf{q}}_1^T, \dots, \bar{\mathbf{q}}_m^T\}^T, \quad \bar{\mathbf{p}} = \{\bar{\mathbf{p}}_1^T, \bar{\mathbf{p}}_2^T, \dots, \bar{\mathbf{p}}_{n-1}^T, \mathbf{p}_n^T\}^T \quad (27)$$

$$\mathbf{M} = \{M_0, M_1, \dots, M_m\}, \bar{\mathbf{M}} = \{M_0, M_1, \dots, M_{m-1}\},$$

$$\bar{\mathbf{M}} = \{M_1, M_2, \dots, M_m\}, \mathbf{N} = \{N_0, N_1, \dots, N_n\},$$

$$\bar{\mathbf{N}} = \{N_0, N_1, \dots, N_{n-1}\}, \bar{\bar{\mathbf{N}}} = \{N_1, N_2, \dots, N_n\} \quad (28)$$

并定义 $k \times l$ 维矩阵 A 和 $s \times t$ 维矩阵 B 的 Kronecker 积为 $ks \times lt$ 维矩阵 $A \otimes B$,即

$$A \otimes B = \begin{bmatrix} a_{11} B & \dots & a_{1l} B \\ \vdots & \ddots & \vdots \\ a_{k1} B & \dots & a_{kl} B \end{bmatrix} \quad (29)$$

那么有

$$\mathbf{q} = (M \otimes I) \bar{\mathbf{q}} \quad \mathbf{p} = (N \otimes I) \{\bar{\mathbf{p}}_0^T \quad \bar{\mathbf{p}}^T\}^T \quad (30)$$

而方程(25)和(26)可表示为

$$\mathbf{F}_1(\bar{\mathbf{q}}, \bar{\mathbf{p}}) = \int_0^n (\bar{\mathbf{M}}^T \otimes \mathbf{p} - \mathbf{M}^T \otimes \frac{\partial H}{\partial \mathbf{q}}) dt + \begin{bmatrix} \mathbf{p}_0 \\ 0 \\ \vdots \\ 0 \\ -\mathbf{p}_n \end{bmatrix} = 0 \quad (31)$$

$$\mathbf{F}_2(\bar{\mathbf{q}}, \bar{\mathbf{p}}) = \int_0^n \bar{\mathbf{N}}^T \otimes (\dot{\bar{\mathbf{q}}} - \frac{\partial H}{\partial \mathbf{p}}) dt + \begin{bmatrix} \bar{\mathbf{q}}_0 - \mathbf{q}_0 \\ 0 \\ \vdots \\ 0 \end{bmatrix} \quad (32)$$

$$\mathbf{q}_m = \bar{\mathbf{q}}_m - \int_0^n \bar{\bar{\mathbf{N}}}_n (\dot{\bar{\mathbf{q}}} - \frac{\partial H}{\partial \mathbf{p}}) dt \quad (33)$$

方程(32)和(33)共包括 $(m+n+1)d$ 个方程,而未知变量 $\bar{\mathbf{q}}$ 和 $\bar{\mathbf{p}}$ 的数量也是 $(m+n+1)d$,因此可由方程(31)和(32)解出 $\bar{\mathbf{q}}$ 和 $\bar{\mathbf{p}}$,然后带入方程(33)求出 \mathbf{q}_m .

一般情况方程(31)和(32)是非线性代数方程组,不能直接解析求解,本文采用牛顿(Newton)迭代法求解,这就需要方程(31)和(32)的 Jacobi 矩阵. 方程(31)和(32)分别对 $\bar{\mathbf{q}}$ 和 $\bar{\mathbf{p}}$ 求导即可得到 Jacobi 矩阵 G ,即

$$G = \begin{bmatrix} G_{11} & G_{12} \\ G_{21} & G_{22} \end{bmatrix} \quad (34)$$

其中

$$G_{11} = \frac{\partial \mathbf{F}_1}{\partial \bar{\mathbf{q}}} = - \int_0^n (\bar{\mathbf{M}}^T \bar{\mathbf{M}}) \otimes \frac{\partial H}{\partial \mathbf{q} \partial \mathbf{q}} dt \quad (35)$$

$$G_{12} = \frac{\partial \mathbf{F}_1}{\partial \bar{\mathbf{p}}} = \int_0^n ((\bar{\mathbf{M}}^T \bar{\mathbf{N}}) \otimes I - (\bar{\mathbf{M}}^T \bar{\mathbf{N}}) \otimes \frac{\partial H}{\partial \mathbf{q} \partial \mathbf{p}}) dt - \begin{bmatrix} 0 & \dots & 0 & 0 \\ \vdots & \ddots & \vdots & \vdots \\ 0 & \dots & 0 & 0 \\ 0 & \dots & 0 & 1 \end{bmatrix}_{(m+1) \times n} \otimes I \quad (36)$$

$$G_{21} = \frac{\partial \mathbf{F}_2}{\partial \bar{\mathbf{q}}} = \int_0^n ((\bar{\mathbf{N}}^T \bar{\mathbf{M}}^T) \otimes I - (\bar{\mathbf{N}}^T \bar{\mathbf{M}}) \otimes \frac{\partial H}{\partial \mathbf{p} \partial \mathbf{q}}) dt - \begin{bmatrix} 1 & \dots & 0 & 0 \\ \vdots & \ddots & \vdots & \vdots \\ 0 & \dots & 0 & 0 \\ 0 & \dots & 0 & 0 \end{bmatrix}_{n \times (m+1)} \otimes I \quad (37)$$

$$G_{22} = \frac{\partial \mathbf{F}_2}{\partial \bar{\mathbf{p}}} = - \int_0^n (\bar{\mathbf{N}}^T \bar{\mathbf{N}}) \otimes \frac{\partial H}{\partial \mathbf{p} \partial \mathbf{p}} dt \quad (38)$$

其中 $\bar{\mathbf{N}}, \bar{\mathbf{N}}, \bar{\bar{\mathbf{N}}}$ 如方程(28). 由于方程(31) - (33) 给出的数值算法基于变分原理,并且完全满足正则变换的条件,因此是自然保辛的数值算法. 一般情况下,方程(31) - (33)和方程(35) - (38)中的积分不可能解析积分,必须借助于数值积分,本文采用 Gauss 积分公式, Gauss 积分点的个数用 g 表示. 这种方法可称为 SPPmng, S 表示保辛(symplectic),第一个 P 表示左端的动量是独立变量,第二个 P 表示右端的动量为独立变量, m 表示位移的近似拉格朗日多项式的阶次, n 表示动量的近似拉格朗日多项式的阶次,而 g 表示 Gauss 积分点的个数.

如果 $n = m + 1$,通过方程(31)和(32)很容易得到

$$\bar{\mathbf{p}} = \left(\int_0^n \bar{\mathbf{M}}^T \bar{\mathbf{N}} dt - \begin{bmatrix} 0 & \dots & 0 & 0 \\ \vdots & \ddots & \vdots & \vdots \\ 0 & \dots & 0 & 0 \\ 0 & \dots & 0 & 1 \end{bmatrix}_{m \times (n-1)} \right)^{-1} \otimes I$$

$$I_d \left(\int_0^\eta M^T \otimes \frac{\partial H}{\partial q} dt - \int_0^\eta \dot{M}^T N_0 dt p_0 - \begin{Bmatrix} p_0 \\ 0 \\ \vdots \\ 0 \end{Bmatrix} \right) \quad (39)$$

$$\bar{q} = \left(\int_0^\eta N^T \dot{M} dt + \begin{bmatrix} 1 & 0 & \cdots & 0 \\ 0 & 0 & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \cdots & 0 \end{bmatrix}_{(n-1) \times m} \right)^{-1} \otimes I_d \left(\int_0^\eta N^T \otimes \frac{\partial H}{\partial p} dt + \begin{Bmatrix} q_0 \\ 0 \\ \vdots \\ 0 \end{Bmatrix} \right) \quad (40)$$

方程(39)和(40)自然形成不动点迭代格式. 不动点迭代格式不需要非线性代数方程的 Jacobi 矩阵, 一般比 Newton 迭代法效率高.

4 数值算例

本文数值算例在 64 位计算机和 64 位操作系统下完成.

表 1 位移和动量采用不同阶次插值所需最少 Gauss 积分点个数

Table 1 The smallest number of Gauss integration points for different order of displacement and momentum

| m = | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 | 9 | 10 |
|-----|---|---|------|------|--------|--------|----------|----------|------------|-------------|
| 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 | 9 | 10 | 11 |
| 2 | 1 | 3 | 4 | 5 | 6 | 7 | 8 | 9 | 10 | 11 |
| 3 | 2 | 2 | 2(3) | 5 | 6 | 7 | 8 | 9 | 10 | 11 |
| 4 | 3 | 3 | 3 | 3(4) | 3(4,5) | 7 | 8 | 9 | 10 | 11 |
| n = | 5 | 4 | 4 | 4 | 4(5) | 4(5,6) | 4(5,6,7) | 9 | 10 | 11 |
| 6 | 5 | 5 | 5 | 5 | 5 | 5(6) | 5(6,7) | 5(6,7,8) | 5(6,7,8,9) | 11 |
| 7 | 6 | 6 | 6 | 6 | 6 | 6 | 6(7) | 6(7,8) | 6(7,8,9) | 6(7,8,9,10) |
| 8 | 7 | 7 | 7 | 7 | 7 | 7 | 7 | 7(8) | 7(8,9) | 7(8,9,10) |
| 9 | 8 | 8 | 8 | 8 | 8 | 8 | 8 | 8 | 8(9) | 8(9,10) |
| 10 | 9 | 9 | 9 | 9 | 9 | 9 | 9 | 9 | 9 | 9(10) |

考虑 SPPmng 算法的精度. 如果算法的精度为 s 阶, 表明如果时间步长为 η , 则算法的误差为 $C\eta^s$, 记为 $O(\eta^s)$. 因此, 如果采用两个不同的时间步长 η_1 和 η_2 积分, 它们积分结果的误差分别为 e_1 和 e_2 , 则有

$$e_1 = C\eta_1^s, \quad e_2 = C\eta_2^s \quad (42)$$

那么算法的精度阶数 s 可通过如下方程计算, 即

$$s = \frac{\log(e_1) - \log(e_2)}{\log(\eta_1) - \log(\eta_2)} \quad (43)$$

本文讨论的是哈密顿系统的积分, 由于哈密顿系统的哈密顿函数是守恒量, 因此, 可以通过数值积分得到的哈密顿函数与真实哈密顿函数来计算算法积分的误差.

考虑单摆方程, 其哈密顿函数为

$$H = \frac{1}{2}p^2 - \cos(q) \quad (41)$$

取初始条件为 $q(0) = 1$ 和 $p(0) = 1$.

对广义位移和动量采用不同阶次拉格朗日多项式插值, 考虑 SPPmng 方法所需最少 Gauss 积分点个数. 表 1 给出了需要最少的 Gauss 积分点个数, 括号中的数字表示不能积分的 Gauss 积分点个数, 例如表 1 中 $m = 8$ 和 $n = 5$ 对应的项为 5(6,7), 这表示取 5 个以及大于或等于 8 个 Gauss 积分点得到的保辛算法存在唯一解, 而取 6 或 7 个 Gauss 积分点得到的保辛算法, 由于在 Newton 迭代求解过程中, Jacobi 矩阵奇异, 因此不能求解. 表 1 表明: 对于 SPPmng 方法, 如果广义位移的插值多项式阶数为 m , 而广义动量的插值多项式阶数为 n , 则若 $m \leq \min(2n - 3, n - 1)$, 本文保辛算法要求 Gauss 积分点个数满足 $g \geq n - 1$; 若 $2n - 3 \geq m \geq n$, 则要求 Gauss 积分点个数满足 $g = n - 1$ 或 $g \geq m + 1$; 若 $m > 2n - 3$, 则要求 Gauss 积分点个数满足 $g \geq m + 1$.

对于本算例的单摆问题, 在 0 到 100(s) 区间上积分, 分别采用 $\eta_1 = 0.4(s)$ 和 $\eta_2 = 0.2(s)$ 两个时间步长, 并计算哈密顿函数的相对误差, 然后利用方程(43)计算算法的精度. 表 2-7 分别给出了位移采用 1 至 6 阶拉格朗日多项式插值, 而动量采用不同阶次拉格朗日多项式插值, 以及采用不同 Gauss 积分点个数时, 方法精度的阶数 s . 综合表 2-7 可以得到如下结论: 对于 SPPmng 方法, 如果广义位移的插值多项式阶数为 m , 而广义动量的插值多项式阶数为 n , 并且 Gauss 积分点个数满足表 1 的条件, 则若 $n \geq m + 2$, 算法精度为 $s = 2m$; 若 $n \leq$

$m + 1$ 且 $g \geq m + 1$, 算法精度为 $s = 2n$; 若 $n \leq m + 1$ 且 $g \leq m$, 算法精度为 $s = 2(n + g - m - 1)$.

表 2 $m = 1$ 而 n 和 g 变化时保辛算法的精度

Table 2 The order of the symplectic method for $m = 1$ and different n and g

| Gauss | n = 1 | n = 2 | n = 3 | n = 4 | n = 5 | n = 6 | n = 7 | n = 8 |
|-------|--------|--------|--------|--------|--------|--------|--------|--------|
| 1 | — | 1.9997 | — | — | — | — | — | — |
| 2 | 2.0069 | 4.0084 | 2.0076 | — | — | — | — | — |
| 3 | 2.0069 | 4.0233 | 2.0082 | 2.0065 | — | — | — | — |
| 4 | 2.0069 | 4.0234 | 2.0096 | 2.0072 | 2.0065 | — | — | — |
| 5 | 2.0069 | 4.0234 | 2.0096 | 2.0081 | 2.0071 | 2.0065 | — | — |
| 6 | 2.0069 | 4.0234 | 2.0096 | 2.0081 | 2.0075 | 2.0068 | 2.0065 | — |
| 7 | 2.0069 | 4.0234 | 2.0096 | 2.0081 | 2.0075 | 2.0071 | 2.0068 | 2.0065 |
| 8 | 2.0069 | 4.0234 | 2.0096 | 2.0081 | 2.0075 | 2.0071 | 2.0070 | 2.0066 |
| 9 | 2.0069 | 4.0234 | 2.0096 | 2.0081 | 2.0075 | 2.0071 | 2.0070 | 2.0068 |

表 3 $m = 2$ 而 n 和 g 变化时保辛算法的精度

Table 3 The order of the symplectic method for $m = 2$ and different n and g

| Gauss | n = 1 | n = 2 | n = 3 | n = 4 | n = 5 | n = 6 | n = 7 | n = 8 |
|-------|--------|--------|--------|--------|--------|--------|--------|--------|
| 1 | — | — | — | — | — | — | — | — |
| 2 | — | — | 4.0084 | — | — | — | — | — |
| 3 | 2.0069 | 4.0239 | 6.0217 | 4.0110 | — | — | — | — |
| 4 | 2.0069 | 4.0249 | 6.0182 | 4.0108 | 4.0104 | — | — | — |
| 5 | 2.0069 | 4.0249 | 6.0182 | 4.0122 | 4.0117 | 4.0104 | — | — |
| 6 | 2.0069 | 4.0249 | 6.0182 | 4.0122 | 4.0119 | 4.0106 | 4.0104 | — |
| 7 | 2.0069 | 4.0249 | 6.0182 | 4.0122 | 4.0119 | 4.0111 | 4.0110 | 4.0104 |
| 8 | 2.0069 | 4.0249 | 6.0182 | 4.0122 | 4.0119 | 4.0111 | 4.0111 | 4.0105 |
| 9 | 2.0069 | 4.0249 | 6.0182 | 4.0122 | 4.0119 | 4.0111 | 4.0111 | 4.0108 |

表 4 $m = 3$ 而 n 和 g 变化时保辛算法的精度

Table 4 The order of the symplectic method for $m = 3$ and different n and g

| Gauss | n = 1 | n = 2 | n = 3 | n = 4 | n = 5 | n = 6 | n = 7 | n = 8 |
|-------|--------|--------|--------|--------|--------|--------|--------|--------|
| 1 | — | — | — | — | — | — | — | — |
| 2 | — | — | 2.0069 | — | — | — | — | — |
| 3 | — | — | — | 6.0217 | — | — | — | — |
| 4 | 2.0069 | 4.0278 | 6.0378 | 8.0434 | 6.0430 | — | — | — |
| 5 | 2.0069 | 4.0278 | 6.0396 | 8.0523 | 6.0396 | 6.0381 | — | — |
| 6 | 2.0069 | 4.0278 | 6.0396 | 8.0528 | 6.0404 | 6.0379 | 6.0381 | — |
| 7 | 2.0069 | 4.0278 | 6.0396 | 8.0528 | 6.0404 | 6.0389 | 6.0387 | 6.0381 |
| 8 | 2.0069 | 4.0278 | 6.0396 | 8.0528 | 6.0404 | 6.0389 | 6.0388 | 6.0381 |
| 9 | 2.0069 | 4.0278 | 6.0396 | 8.0528 | 6.0404 | 6.0389 | 6.0388 | 6.0385 |

表 5 $m = 4$ 而 n 和 g 变化时保辛算法的精度

Table 5 The order of the symplectic method for $m = 4$ and different n and g

| Gauss | n = 1 | n = 2 | n = 3 | n = 4 | n = 5 | n = 6 | n = 7 | n = 8 |
|-------|--------|--------|--------|--------|---------|--------|--------|--------|
| 1 | — | — | — | — | — | — | — | — |
| 2 | — | — | — | — | — | — | — | — |
| 3 | — | — | — | 4.0239 | — | — | — | — |
| 4 | — | — | — | — | 8.0434 | — | — | — |
| 5 | 2.0069 | 4.0279 | 6.0367 | 8.0301 | 10.0478 | 8.0281 | — | — |
| 6 | 2.0069 | 4.0279 | 6.0367 | 7.9902 | 10.0352 | 8.0265 | 8.0264 | — |
| 7 | 2.0069 | 4.0279 | 6.0367 | 7.9904 | 10.0355 | 8.0278 | 8.0275 | 8.0264 |
| 8 | 2.0069 | 4.0279 | 6.0367 | 7.9904 | 10.0355 | 8.0278 | 8.0276 | 8.0264 |
| 9 | 2.0069 | 4.0279 | 6.0367 | 7.9904 | 10.0355 | 8.0278 | 8.0276 | 8.0270 |

表6 $m=5$ 而 n 和 g 变化时保辛算法的精度Table 6 The order of the symplectic method for $m=5$ and different n and g

| Gauss | $n=1$ | $n=2$ | $n=3$ | $n=4$ | $n=5$ | $n=6$ | $n=7$ | $n=8$ |
|-------|--------|--------|--------|--------|---------|---------|---------|---------|
| 1 | — | — | — | — | — | — | — | — |
| 2 | — | — | — | — | — | — | — | — |
| 3 | — | — | — | 2.0069 | — | — | — | — |
| 4 | — | — | — | — | 6.0378 | — | — | — |
| 5 | — | — | — | — | — | 10.0478 | — | — |
| 6 | 2.0069 | 4.0279 | 6.0367 | 7.9808 | 10.1020 | 12.0882 | 10.0670 | — |
| 7 | 2.0069 | 4.0279 | 6.0367 | 7.9813 | 10.1007 | 12.0673 | 10.0638 | 10.0633 |
| 8 | 2.0069 | 4.0279 | 6.0367 | 7.9813 | 10.1007 | 12.0682 | 10.0654 | 10.0646 |
| 9 | 2.0069 | 4.0279 | 6.0367 | 7.9813 | 10.1007 | 12.0682 | 10.0654 | 10.0650 |

表7 $m=6$ 而 n 和 g 变化时保辛算法的精度Table 7 The order of the symplectic method for $m=6$ and different n and g

| Gauss | $n=1$ | $n=2$ | $n=3$ | $n=4$ | $n=5$ | $n=6$ | $n=7$ | $n=8$ |
|-------|--------|--------|--------|--------|---------|---------|---------|---------|
| 1 | — | — | — | — | — | — | — | — |
| 2 | — | — | — | — | — | — | — | — |
| 3 | — | — | — | — | — | — | — | — |
| 4 | — | — | — | — | 4.0278 | — | — | — |
| 5 | — | — | — | — | — | 8.0301 | — | — |
| 6 | — | — | — | — | — | — | 12.0882 | — |
| 7 | 2.0069 | 4.0279 | 6.0367 | 7.9829 | 10.1052 | 12.0837 | 14.1331 | 12.0591 |
| 8 | 2.0069 | 4.0279 | 6.0367 | 7.9828 | 10.1052 | 12.0842 | 14.0813 | 12.0565 |
| 9 | 2.0069 | 4.0279 | 6.0367 | 7.9829 | 10.1052 | 12.0842 | 14.0823 | 12.0580 |

表8 $m=5$ 而 n 和 g 变化时单摆哈密顿函数的相对误差Table 8 The relative error of the Hamilton function for $m=5$ and different n and g

| Gauss | $n=1$ | $n=2$ | $n=3$ | $n=4$ | $n=5$ | $n=6$ | $n=7$ | $n=8$ |
|-------|-----------|-----------|-----------|------------|------------|------------|------------|------------|
| 1 | — | — | — | — | — | — | — | — |
| 2 | — | — | — | — | — | — | — | — |
| 3 | — | — | — | $5.74e-2$ | — | — | — | — |
| 4 | — | — | — | — | $1.08e-7$ | — | — | — |
| 5 | — | — | — | — | — | $5.60e-14$ | — | — |
| 6 | $5.74e-2$ | $8.80e-5$ | $1.08e-7$ | $2.31e-10$ | $9.60e-14$ | $4.66e-17$ | $9.16e-15$ | — |
| 7 | $5.74e-2$ | $8.80e-5$ | $1.08e-7$ | $2.31e-10$ | $9.60e-14$ | $1.96e-17$ | $9.13e-15$ | $9.13e-15$ |
| 8 | $5.74e-2$ | $8.80e-5$ | $1.08e-7$ | $2.31e-10$ | $9.60e-14$ | $1.96e-17$ | $5.28e-15$ | $5.27e-15$ |
| 9 | $5.74e-2$ | $8.80e-5$ | $1.08e-7$ | $2.31e-10$ | $9.60e-14$ | $1.96e-17$ | $5.28e-15$ | $5.27e-15$ |

表8具体给出了时间步长为 $0.1(s)$ 时, $m=5$ 而 n 和 g 变化时,单摆哈密顿函数的相对误差.表8表明:对于SPPmng方法,对于给定 m ,最优的组合为 $n=m+1$ 和 $g=m+2$,而这种组合恰好可导出方程(39)和(40)表示的不动点格式,迭代求解效率更高.

5 结论

本文将广义位移和动量同时用拉格朗日多项式近似,并选择积分区间两端的广义动量为独立变量,然后基于对偶变量变分原理导出了哈密顿系统

的离散正则变换和对应的数值积分保辛算法.通过数值算例得到如下结论:

(1) 如果广义位移的插值多项式阶数为 m ,而广义动量的插值多项式阶数为 n ,Gauss积分点个数为 g ,则若 $m \leq \min(2n-1, n+1)$,本文保辛算法要求Gauss积分点个数满足 $g \geq n-1$;若 $2n-3 \geq m \geq n$,则要求Gauss积分点个数满足 $g = n-1$ 或 $g \geq m+1$;若 $m > 2n-3$,则要求Gauss积分点个数满足 $g \geq m+1$.

(2) 在Gauss积分点个数满足结论1的条件下,若 $n \geq m+2$,算法精度为 $s=2m$;若 $n \leq m+1$ 且

$g \geq m + 1$, 算法精度为 $s = 2n$; 若 $n \leq m + 1$ 且 $g \leq m$, 算法精度为 $s = 2(n + g - m - 1)$.

(3) 对于给定 m , 最优的组合为 $n = m + 1$ 和 $g = m + 2$, 而这种组合恰好可导出不动点格式, 迭代求解效率更高.

参 考 文 献

- 1 Arnold V I. *Mathematical Methods of Classical Mechanics*. New York: Springer - Verlag, 1989
- 2 Goldstein H. *Classical Mechanics*. 2 ed. London: Addison - Wesley, 1980
- 3 冯康, 秦孟兆. 哈密尔顿系统的辛几何算法. 杭州: 浙江科学技术出版社, 2003 (Feng Kang, Qin Mengzhao. *Symplectic Geometric Algorithm for Hamiltonian Systems*. Hangzhou: Zhejiang Science and Technology Press, 2003 (in Chinese))
- 4 Feng K. On Difference Schemes and Symplectic Geometry. *Proceedings of the 5th international symposium on differential geometry and differential equations*, Beijing, 1984
- 5 Hairer E, Wanner G. *Solving Ordinary Differential Equations II - Stiff and Differential - Algebraic Problems* 2ed. Berlin: Springer, 1996
- 6 Hairer E, Nørsett S P, Wanner G. *Solving Ordinary Differential Equations I - Nonstiff Problems* 2ed. Berlin: Springer, 1993
- 7 Hairer E, Lubich C, Wanner G. *Geometric Numerical Integration: Structure - Preserving Algorithm for Ordinary Differential Equations*. Second ed. New York: Springer, 2006
- 8 钟万勰. 分析结构力学与有限元. 动力学与控制学报, 2004, 2(4): 1 ~ 8 (Zhong Wanxie. *Analytical Structural Mechanics and Finite Element*. *Journal of Dynamics and Control*, 2004, 2(4): 1 ~ 8 (in Chinese))
- 9 钟万勰, 姚征. 时间有限元与保辛. 机械强度, 2005, 27(2): 178 ~ 183 (Zhong Wanxie, Yao Zheng. *Time Domain FEM and Symplectic Conservation*. *Journal of Mechanical Strength*, 2005, 27(2): 178 ~ 183 (in Chinese))
- 10 钟万勰, 高强. 约束动力系统的分析结构力学积分. 动力学与控制学报, 2006, 4(3): 193 ~ 200 (Zhong Wanxie, Gao Qiang. *Integration of Constrained Dynamical System Via Analytical Structural Mechanics*. *Journal of Dynamics and Control*, 2006, 4(3): 193 ~ 200 (in Chinese))

SYMPLECTIC METHOD BASED ON DUAL VARIABLE PRINCIPLE AND INDEPENDENT MOMENTUM AT TWO ENDS *

Gao Qiang Tan Shujun Zhang Hongwu Zhong Wanxie

(*Department of Engineering Mechanics, State Key Laboratory of Structural Analysis of Industrial Equipment, Dalian University of Technology, Dalian 116023*)

Abstract The generalized displacements and momentum were approximated by Lagrange polynomial, and the momentum at the two ends of time interval were taken as the independent variables. Then the discrete Hamilton canonical equations and the corresponding symplectic method were derived based on the dual variable principle. A fix point iteration formula can be derived when the order of the approximate polynomials of displacements and momentum satisfy some certain conditions. In the numerical examples part, the smallest number of Gauss integration point required for different order of the approximate polynomials of displacements and momentum was discussed, and also the numerical precision of the proposed symplectic method for different orders of the approximate polynomials of displacements and momentum and numbers of Gauss integration point was discussed. The fix point iteration formula is the optimal one.

Key words variable principle, symplectic method, Hamilton system, dual

Received 10 December 2008, revised 30 December 2008.

* The project supported by the National Natural Science Foundation of China (10632030, 10721062, 2005CB321704), the doctoral research fund of Liaoning (20081091), Young Researcher Funds of Dalian University of Technology and Science Research Foundation of Dalian University of Technology (SFDUT07002)